



HAL
open science

Well-posedness of an interaction model on Riemannian manifolds

Razvan C. Fetecau, Francesco S. Patacchini

► **To cite this version:**

Razvan C. Fetecau, Francesco S. Patacchini. Well-posedness of an interaction model on Riemannian manifolds. *Communications on Pure and Applied Analysis*, 2022, 21 (11), pp.3559-3585. 10.3934/cpaa.2022114 . hal-03877707v1

HAL Id: hal-03877707

<https://ifp.hal.science/hal-03877707v1>

Submitted on 6 Dec 2022 (v1), last revised 30 Mar 2023 (v2)

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.



Distributed under a Creative Commons Attribution - NonCommercial - NoDerivatives 4.0 International License



WELL-POSEDNESS OF AN INTERACTION MODEL ON RIEMANNIAN MANIFOLDS

RAZVAN C. FETEAU

Department of Mathematics
Simon Fraser University
Burnaby, BC V5A 1S6, Canada

FRANCESCO S. PATACCINI*

IFP Energies nouvelles
1-4 Avenue de Bois-Préau
92852 Rueil-Malmaison, France

(Communicated by Dong Ye)

ABSTRACT. We investigate a model for collective behaviour with intrinsic interactions on smooth Riemannian manifolds. For regular interaction potentials, we establish the local well-posedness of measure-valued solutions defined via optimal mass transport. We also extend our result to the global well-posedness of solutions for manifolds with nonpositive bounded sectional curvature. The core concept underlying the proofs is that of Lipschitz continuous vector fields in the sense of parallel transport.

1. Introduction. This paper is concerned with the following integro-differential equation for the evolution of a population density ρ on a Riemannian manifold M :

$$\partial_t \rho - \nabla_M \cdot (\rho \nabla_M K * \rho) = 0, \quad (1.1)$$

where $K: M \times M \rightarrow \mathbb{R}$ is an interaction (also known as aggregation) potential which models social interactions such as attraction and repulsion, and $\nabla_M \cdot$ and ∇_M represent the manifold divergence and gradient, respectively. In (1.1) the symbol $*$ denotes a measure convolution: for a time-dependent measure ρ_t on M and $x \in M$ we set

$$K * \rho_t(x) = \int_M K(x, y) d\rho_t(y).$$

We restrict our solutions ρ to be probability measures on M at all times: $\int_M d\rho_t = 1$ for all t .

Model (1.1) has numerous applications in swarming and self-organized behaviour in biology [34], material science [10], robotics [22, 28], and social sciences [35]. In such applications, equation (1.1) can model interactions between biological organisms such as insects, birds or cells, as well as interactions between robots or even opinions. Concerning the theory, the mathematical analysis of solutions to model (1.1) has

2020 *Mathematics Subject Classification.* Primary: 35A01, 35B40; Secondary: 37C05, 58J90.

Key words and phrases. Intrinsic interactions, measure solutions, swarming on manifolds, parallel transport, Lipschitz continuity.

The first author is supported by an NSERC Canada Discovery grant.

* Corresponding author.

focused almost exclusively on the model set up on Euclidean space \mathbb{R}^n ; we refer in this case to [5–8] for the well-posedness of the initial-value problem and to [4, 15, 18, 19, 32] for the long-time behaviour of solutions.

The goal of the paper is to establish the well-posedness of measure-valued solutions to model (1.1) set up on general Riemannian manifolds. There are only very few works on this subject, and all require that the manifold is embedded in a larger Euclidean space. For instance, in [9, 36, 38], the authors investigate the well-posedness of the aggregation model (1.1) when the interactions are *extrinsic*, in the sense that the interaction potential depends on the embedding Euclidean distance between points. Another class of models considers *intrinsic* interactions, where the interaction potential depends on the geodesic distance on M between points. Recent works in this direction investigate the well-posedness of model (1.1) with intrinsic interactions on the sphere [20] and on the special orthogonal group $SO(3)$ [17]. In both [17, 20], however, certain calculations make use of the ambient vector spaces of the manifolds, i.e., in \mathbb{R}^n and $\mathbb{R}^{3 \times 3}$, respectively.

This paper presents a fully intrinsic approach to the well-posedness of solutions to (1.1) on Riemannian manifolds, which does not require any extrinsic calculations in an ambient vector space. To this aim we use the concept of Lipschitz continuity of vector fields via parallel transport. While this is a concept widely used in the literature on optimization on manifolds [12, 16], it is much less common in the analysis of differential equations on manifolds. Lipschitz continuity by parallel transport enables us to compare tangent vectors in an intrinsic manner. This approach fundamentally distinguishes itself from that in [17, 20], where tangent vectors at different points on the manifold are compared in the norm of an ambient vector space.

As a further motivation for this paper we also mention various studies on the long-time behaviour of solutions to (1.1) on manifolds. An interesting collection of equilibria on the sphere and the hyperbolic plane for model (1.1) with intrinsic interactions can be found in [21]. These equilibria show a very rich pattern formation behaviour (e.g., disks, annuli, rings), similar to what has been observed in \mathbb{R}^n [30, 37]. In addition, emergent behaviour has been studied extensively in the related Lohe-type models with extrinsic interactions on the unit sphere, matrix manifolds and tensor spaces in [23–25]. In these works the focus is to investigate the formation of consensus solutions, where the equilibria consist of an aggregation at a single point, rather than the well-posedness of (1.1) in the intrinsic setting.

The paper is structured as follows. We first review in Section 2 the concept of parallel transport on a Riemannian manifold and that of Lipschitz continuity by parallel transport, and present useful relations in terms of Hessians and flow maps. In Section 3 we give the rigorous notion of solution to equation (1.1) which will be used in the well-posedness theory, as well as introduce the 1-Wasserstein probability space, which determines the regularity of the solutions we seek. Then, in Section 4 we present our proof of local and global well-posedness for equation (1.1), starting with the fundamental lemmas underlying the proof. In this section, our results on the global well-posedness are restricted to manifolds with nonpositive bounded sectional curvature. Finally, in Section 5 we state two additional results which can be obtained by simply adapting results already proven for the sphere in [20]. Appendices A and B give, respectively, basic notions of differential geometry useful to our purposes, and the proofs to preliminary results stated in the main body of the paper.

Everywhere in this paper, M denotes a manifold satisfying the assumption that follows. We shall only refer to it in some of the main statements; anywhere else, it will be implicitly assumed.

(M) M is a complete, smooth Riemannian manifold of finite dimension n , with positive injectivity radius. We denote its intrinsic distance by d and its sectional curvature by \mathcal{K} .

2. Preliminaries on parallel transport and Lipschitz continuity. In Appendix A we review some basic concepts and terminology from differential geometry that are relevant to the present work, and also introduce some notation. We invite the reader to look there whenever unfamiliar with some of the concepts or notation used in the main body of the paper. In particular, in Appendix A we briefly discuss the logarithm map, normal and totally normal neighbourhoods, the injectivity radius, geodesics, convex sets, normal charts and the push-forward, as well as recall some of their properties important to our analysis; we also give a useful relationship between the Euclidean norm and the intrinsic distance on M .

Anywhere in the paper, for $x \in M$ we write $\langle \cdot, \cdot \rangle_x$ and $\|\cdot\|_x$ the tangent inner product and norm, respectively, on T_xM , the tangent space of M at x . The tangent bundle of M is denoted by TM .

In this section, we present some background on Lipschitz continuity by parallel transport and its link to Hessians, as well as give the Cauchy–Lipschitz theorem for flow maps on manifolds.

2.1. Lipschitz continuity via parallel transport. Given a curve $\gamma: [0, 1] \rightarrow M$ and $v \in T_{\gamma(0)}M$, the *parallel transport* of v along γ is given by the unique solution $X: [0, 1] \rightarrow TM$ with $X(t) \in T_{\gamma(t)}M$ for all $t \in [0, 1]$ to the ODE

$$\begin{cases} \nabla_{\gamma'(t)}X(t) = 0, \\ X(0) = v, \end{cases}$$

where $\nabla_{\gamma'}X$ denotes the covariant derivative of X along γ and ∇ is the Levi-Civita connection on M . We denote this solution by $\Pi_{\gamma,t}v$ for all $t \in [0, 1]$. Given $t \in [0, 1]$, the map $v \mapsto \Pi_{\gamma,t}v$ is a linear isometry from $T_{\gamma(0)}M$ to $T_{\gamma(t)}M$, i.e.,

$$\langle \Pi_{\gamma,t}v, \Pi_{\gamma,t}w \rangle_{\gamma(t)} = \langle v, w \rangle_{\gamma(0)}, \quad \text{for all } v, w \in T_{\gamma(0)}M.$$

The situation of interest in this paper is when $\gamma: [0, 1] \rightarrow M$ is the *unique* minimizing geodesic connecting a point x to a point y . In this case we write Π_{xy} for $\Pi_{\gamma,1}$. It holds that $\Pi_{xy}^{-1} = \Pi_{yx}$, which means that taking the parallel transport of a vector $v \in T_xM$ to T_yM and then back to T_xM returns the same vector v .

A *vector field* on a set $U \subset M$ is a function $X: U \rightarrow TM$ such that $X(x) \in T_xM$ for all $x \in U$. Using parallel transport, we can give the following definition of Lipschitz continuous vector fields. We use in our setup a totally normal neighbourhood of M , as any two points in such a set can be connected by a unique minimizing geodesic (though not necessarily lying entirely in it).

Definition 2.1 (Lipschitz continuity via parallel transport). Suppose that U is a totally normal neighbourhood of M and let X be a vector field on U . We say that X is *locally Lipschitz continuous* if for all compact sets $Q \subset U$ there exists a constant $L_Q > 0$ such that

$$\|X(x) - \Pi_{yx}X(y)\|_x \leq L_Q d(x, y), \quad \text{for all } x, y \in Q; \tag{2.1}$$

we write $\|X\|_{\text{Lip}(Q)}$ for the smallest such constant. We say that X is (*globally*) *Lipschitz continuous* if there exists $L > 0$ such that (2.1) holds for all $x, y \in U$ by replacing L_Q with L ; we write $\|X\|_{\text{Lip}(U)}$ for the smallest such L .

We note that a locally Lipschitz continuous vector field X on a totally normal neighbourhood U is also locally bounded in the sense that for all $Q \subset U$ compact there exists a constant $C_Q > 0$ such that $\|X(x)\|_x \leq C_Q$ for all $x \in Q$; we denote by $\|X\|_{L^\infty(Q)}$ the smallest such constant.

We also note that this definition of Lipschitz continuity is chart-free. The notion of Lipschitz continuity on charts is the standard concept used to show the well-posedness of flow maps via the Cauchy–Lipschitz theorem in the theory of dynamical systems on manifolds [20, 27]. For completeness, we recall the definition here: a vector field X on an open subset U of M is *locally Lipschitz continuous on charts* if for every chart (V, φ) of M and compact set $Q \subset U \cap V$, there exists a constant $L_{\varphi, Q} > 0$ such that

$$\|\varphi_*X(x) - \varphi_*X(y)\|_{\mathbb{R}^n} \leq L_{\varphi, Q} \|\varphi(x) - \varphi(y)\|_{\mathbb{R}^n}, \quad \text{for all } x, y \in Q. \quad (2.2)$$

The link between local Lipschitz continuity in the sense of Definition 2.1 and that on charts is given in the following lemma, which, in short, states that Lipschitz continuity in the sense of parallel transport implies Lipschitz continuity on *normal* charts.

Lemma 2.2. *Let $x \in M$ and set $\delta \leq r_{\text{conv}}(x)$. Suppose that X is a vector field on the geodesic ball $B_\delta(x)$ which is locally Lipschitz continuous. Let moreover (V, φ) be a normal chart generated by x . Then, for any $Q \subset B_\delta(x) \cap V$ compact there exists $L_{\varphi, Q} > 0$ such that (2.2) holds; furthermore, the constant $L_{\varphi, Q}$ depends linearly on $\|X\|_{L^\infty(Q)}$ and $\|X\|_{\text{Lip}(Q)}$.*

Proof. See Appendix B.1. □

2.2. Hessians. Given a vector field X on an open subset U of M , a point $x \in U$ and a nonzero tangent vector $v \in T_xM$, the *derivative of X at x in direction v* is given, whenever the limit below exists, by

$$\nabla_v X(x) = \left. \frac{d}{dt} \right|_{t=0} \Pi_{x\gamma(t)}^{-1} X(\gamma(t)) = \lim_{t \rightarrow 0} \frac{\Pi_{x\gamma(t)}^{-1} X(\gamma(t)) - X(x)}{t}, \quad (2.3)$$

where γ denotes the geodesic starting at x with velocity v defined at any t small enough for $\exp_x(tv)$ to exist. When indeed the limit above exists, we have $\nabla_v X(x) \in T_xM$ and we say that X is *differentiable at x in direction v* . If X is differentiable at x in every direction, the map $\nabla X(x): T_xM \rightarrow T_xM$ is linear and is called the *derivative of X at x* . If X is differentiable at all points in U in every direction, then we simply say that X is *differentiable* and the map $U \ni x \mapsto \nabla X(x)$ is simply referred to as the *derivative of X* .

For a differentiable function $f: U \rightarrow \mathbb{R}$ such that the vector field ∇f is differentiable at x in direction v , we write $\text{Hess}_v f(x)$ for the *Hessian of f at x in direction v* , defined by

$$\text{Hess}_v f(x) = \nabla_v(\nabla f)(x).$$

If ∇f is differentiable at x in every direction, the map $\text{Hess } f(x): T_xM \rightarrow T_xM$ is linear and is referred to as the *Hessian operator of f at x* . If furthermore ∇f is differentiable at all points in U in every direction, then we simply say that f has a *Hessian* and the map $U \ni x \mapsto \text{Hess } f(x)$ is called the *Hessian of f* . Similar to the

Euclidean case, the lemma below links bounded Hessians to Lipschitz continuity in the parallel transport sense [16]:

Lemma 2.3. *Let $U \subset M$ be a totally normal neighbourhood and let $f : U \rightarrow \mathbb{R}$ have a Hessian. If the gradient of f is locally Lipschitz continuous, then its Hessian is locally bounded, i.e., for all $Q \subset U$ compact there is $C_Q > 0$ such that*

$$\|\text{Hess}_v f(x)\|_x \leq C_Q \|v\|_x, \quad \text{for all } x \in Q \text{ and } v \in T_x M.$$

If furthermore U is geodesically convex, then the converse is also true.

Proof. See Appendix B.2. □

We directly have the following global version of Lemma 2.3:

Lemma 2.4. *Let $U \subset M$ be a totally normal neighbourhood and let $f : U \rightarrow \mathbb{R}$ have a Hessian. If the gradient of f is Lipschitz continuous, then its Hessian is bounded, i.e., there is $C > 0$ such that*

$$\|\text{Hess}_v f(x)\|_x \leq C \|v\|_x, \quad \text{for all } x \in U \text{ and } v \in T_x M.$$

If furthermore U is geodesically convex, then the converse is also true.

In the optimization literature, the property of a function f defined on a totally normal neighbourhood $U \subset M$ to have a Lipschitz continuous gradient, is sometimes referred to as L -smoothness [1, 2, 26]. Lemma 2.4 therefore says that L -smoothness and Hessian boundedness are equivalent on geodesically convex subsets of M (as they are in Euclidean spaces).

2.3. Flow maps and the Cauchy–Lipschitz theorem. A time-dependent vector field on a subset $U \times [0, T)$ of $M \times [0, \infty)$ is a function $X : U \times [0, T) \rightarrow TM$ such that $X(\cdot, t)$ is a vector field on U for all $t \in [0, T)$; we will also use X_t to denote $X(\cdot, t)$. Given such a time-dependent vector field X and $\Sigma \subset U$ measurable, a flow map generated by (X, Σ) is a function $\Psi_X : \Sigma \times [0, \tau) \rightarrow U$, for some $\tau \leq T$, that for all $x \in \Sigma$ and $t \in [0, \tau)$ satisfies

$$\begin{cases} \frac{d}{dt} \Psi_X^t(x) = X_t(\Psi_X^t(x)), \\ \Psi_X^0(x) = x, \end{cases} \tag{2.4}$$

where we used the abbreviation Ψ_X^t for $\Psi_X(\cdot, t)$. A flow map is said to be maximal if its time domain cannot be extended while (2.4) holds; it is said to be global if $\tau = T = \infty$ and local otherwise.

The following theorem gives the local well-posedness of flow maps generated by Lipschitz vector fields in the sense of Definition 2.1.

Theorem 2.5 (Local Cauchy–Lipschitz Theorem). *Suppose that U is a totally normal neighbourhood of M . Let $T \in (0, \infty]$ and let X be a time-dependent vector field on $U \times [0, T)$. Suppose that the vector fields in $\{X_t\}_{t \in [0, T)}$ are locally Lipschitz continuous and satisfy, for any compact sets $Q \subset U$ and $S \subset [0, T)$,*

$$\int_S \left(\|X_t\|_{L^\infty(Q)} + \|X_t\|_{\text{Lip}(Q)} \right) dt < \infty. \tag{2.5}$$

Then, for every compact subset Σ of U , there exists a unique maximal flow map generated by (X, Σ) .

Proof. The proof follows very closely that given in [20] of the more classical version of the theorem for Lipschitz continuous vector fields on charts, with the additional use of Lemma 2.2. For the details, see Appendix B.3. \square

Theorem 2.5 gives *local* well-posedness for flow maps. As given by the theorem below, we have *global* well-posedness whenever M is convex and the vector fields in question are globally Lipschitz on $M \times [0, \infty)$.

Theorem 2.6 (Global Cauchy–Lipschitz Theorem on Convex Manifold). *Assume M is geodesically convex and let X be a time-dependent vector field on $M \times [0, \infty)$. Suppose that the vector fields in $\{X_t\}_{t \in [0, \infty)}$ are Lipschitz continuous and for any compact set $Q \subset M$ there holds*

$$\sup_{t \in [0, \infty)} \left(\|X_t\|_{L^\infty(Q)} + \|X_t\|_{\text{Lip}(M)} \right) < \infty.$$

Then, for every compact subset Σ of M , there exists a unique global flow map generated by (X, Σ) .

Proof. The proof follows from the Escape Lemma if M is compact and from a classical extension argument if M is unbounded. We refer the reader to Appendix B.4 for the details. \square

3. Preliminaries on the interaction equation. For simplicity, as we have already implicitly done in the previous section, we will drop the subindices M on the differential operators in equation (1.1). We present in this section the notion of measure-valued solutions for (1.1) and the 1-Wasserstein probability space. We only give the minimal tools we will need for our purpose; for a thorough theory of probability spaces and continuity equations in Euclidean space, we refer the reader to [3].

3.1. Notion of solution for the interaction equation. For $U \subset M$ open, denote by $\mathcal{P}(U)$ the set of Borel probability measures on the metric space (U, d) and by $C([0, T]; \mathcal{P}(U))$ the set of continuous curves from $[0, T]$ into $\mathcal{P}(U)$ endowed with the narrow topology (i.e., the topology dual to the space of continuous bounded functions on U ; see [3]).

If $\Psi: \Sigma \rightarrow U$ for some measurable $\Sigma \subset U$, we denote by $\Psi\#\rho$ the *push-forward* in the mass transport sense of ρ through Ψ . Equivalently, $\Psi\#\rho$ is the probability measure such that for every measurable function $\zeta: U \rightarrow [-\infty, \infty]$ with $\zeta \circ \Psi$ integrable with respect to ρ , we have

$$\int_U \zeta(x) \, d(\Psi\#\rho)(x) = \int_\Sigma \zeta(\Psi(x)) \, d\rho(x).$$

For $T \in (0, \infty]$ and a curve $(\rho_t)_{t \in [0, T]} \subset \mathcal{P}(U)$, we denote by $v[\rho]: U \times [0, T] \rightarrow TM$ the velocity vector field associated to (1.1), that is,

$$v[\rho](x, t) = -\nabla K * \rho_t(x), \quad \text{for all } (x, t) \in U \times [0, T], \quad (3.1)$$

where for convenience we used ρ_t in place of $\rho(t)$, as we shall often do in the sequel. Whether $v[\rho]$ is well-defined or not depends on U ; indeed, we shall see that for ∇K to make sense in the convolution with ρ_t , the set U should be such that $\log_x y$ exists for all $x, y \in U$.

We define solutions to (1.1) in a geometric way, as the push-forward of the initial data through the corresponding flow map [3, Chapter 8.1]. Specifically, we adopt the following definition as solution to equation (1.1):

Definition 3.1 (Solution). Given $U \subset M$ open, we say that $(\rho_t)_{t \in [0, T]} \subset \mathcal{P}(U)$ is a *weak solution* to (1.1) if $(v[\rho], \text{supp}(\rho_0))$ generates a unique flow map $\Psi_{v[\rho]}$ defined on $\text{supp}(\rho_0) \times [0, T]$, and ρ_t satisfies the implicit relation

$$\rho_t = \Psi_{v[\rho]}^t \# \rho_0, \quad \text{for all } t \in [0, T].$$

It can be shown that solutions in the sense of Definition 3.1 are also weak solutions in the sense of distributions to equation (1.1); see [3, Lemma 8.1.6].

3.2. Wasserstein distance. We will use the *intrinsic 1-Wasserstein distance* to compare solutions to (1.1). Let $U \subset M$ be open. For $\rho, \sigma \in \mathcal{P}(U)$, this distance is defined as:

$$W_1(\rho, \sigma) = \inf_{\pi \in \Pi(\rho, \sigma)} \int_{U \times U} d(x, y) \, d\pi(x, y),$$

where $\Pi(\rho, \sigma) \subset \mathcal{P}(U \times U)$ is the set of transport plans between ρ and σ , i.e., the set of elements in $\mathcal{P}(U \times U)$ with first and second marginals ρ and σ , respectively.

Denote by $\mathcal{P}_1(U)$ the set of probability measures on U with finite first moment and by $\mathcal{P}_\infty(U) \subset \mathcal{P}_1(U)$ the set of probability measures on U with compact support. Both spaces $(\mathcal{P}_1(U), W_1)$ and $(\mathcal{P}_\infty(U), W_1)$ are well-defined metric spaces. In addition, we metrize the space $C([0, T]; \mathcal{P}_1(U))$ (and thus $C([0, T]; \mathcal{P}_\infty(U))$) with the distance \mathbf{W}_1 defined by

$$\mathbf{W}_1(\rho, \sigma) = \sup_{t \in [0, T]} W_1(\rho_t, \sigma_t), \quad \text{for all } \rho, \sigma \in C([0, T]; \mathcal{P}_1(U)).$$

Of course, when U is bounded we have $\mathcal{P}(U) = \mathcal{P}_1(U) = \mathcal{P}_\infty(U)$.

The lemma below, used later in the paper, contains various Lipschitz properties involving the distance W_1 .

Lemma 3.2. *Let $U \subset M$ be open.*

- (i) *Let $\Sigma \subset U$. Let furthermore $\rho \in \mathcal{P}_1(U)$ with $\text{supp}(\rho) \subset \Sigma$ and $\Psi_1, \Psi_2: \Sigma \rightarrow U$ be measurable functions. Then,*

$$W_1(\Psi_1 \# \rho, \Psi_2 \# \rho) \leq \sup_{x \in \text{supp}(\rho)} d(\Psi_1(x), \Psi_2(x)).$$

- (ii) *Let $T \in (0, \infty]$ and let X be a time-dependent vector field on $U \times [0, T]$. Let $\rho \in \mathcal{P}_1(U)$ and suppose that $(X, \text{supp}(\rho))$ generates a flow map Ψ_X defined on $\text{supp}(\rho) \times [0, \tau]$ for some $\tau \leq T$. Suppose furthermore that X is bounded, i.e., there exists $C > 0$ such that $\|X(x, t)\|_{x \in U} < C$ for all $x \in U$ and $t \in [0, T]$. Then,*

$$W_1(\Psi_X^t \# \rho, \Psi_X^s \# \rho) \leq C|t - s|, \quad \text{for all } t, s \in [0, \tau].$$

Proof. We refer to [20, Lemma 2.3] for the proof of this result. □

4. Well-posedness of the interaction equation. We study the well-posedness of solutions to (1.1) with an interaction potential K that satisfies:

(K) $K: M \times M \rightarrow \mathbb{R}$ has the form

$$K(x, y) = g(d(x, y)^2), \quad \text{for all } x, y \in M, \tag{4.1}$$

where $g: [0, \infty) \rightarrow \mathbb{R}$ is differentiable.

Because the interaction equation (1.1) involves the gradient of K , we are interested in the Lipschitz continuity of the gradient of the squared distance function d . For all $x, y \in M$ define $d_y(x) = d(x, y)$; if x and y have a unique minimizing geodesic linking them, it holds that

$$\log_x y = -\frac{1}{2}\nabla d_y^2(x). \quad (4.2)$$

Similarly, for all $x, y \in M$ write $K_y(x) = K(x, y)$, and again if there exists a unique minimizing geodesic linking x and y the chain rule yields

$$\nabla K_y(x) = -2g'(d(x, y)^2)\log_x y. \quad (4.3)$$

In other words, equations (4.2) and (4.3) only hold for points y away from the cut locus of x . In particular, they hold for points $x, y \in M$ with $d(x, y) < \text{inj}(M)$.

Equation (1.1) can be interpreted as an aggregation model using (3.1) and (4.3). Indeed, when a point mass at x interacts with a point mass at y , the mass at x is driven by a force of magnitude proportional to $|g'(d(x, y))^2|d(x, y)$, to move either towards y (provided $g'(d(x, y)^2) > 0$) or away from y (provided $g'(d(x, y)^2) < 0$). The velocity field at x computed by (3.1) accounts for all contributions from interactions with point masses $y \in U$ through the convolution.

4.1. Fundamental lemmas. We give here several fundamental lemmas which will be at the core of our proof of the well-posedness of solutions to (1.1). Note first that, given $x, y \in M$ and $U \subset M$ a normal neighbourhood of both x and y , there holds

$$\Pi_{yx}\log_y(x) = -\log_x(y). \quad (4.4)$$

Equation (4.4) follows from the definition of Riemannian logarithm and the fact that the parallel transport is an isometry [13]. Indeed, the vector $\log_y(x) \in T_yM$ is tangent to the geodesic γ joining y and x , pointing in the direction of x , and has length $\|\log_y(x)\|_y = d(x, y)$. By definition of a geodesic, taking the parallel transport of $\log_y(x)$ from y to x along γ results into a vector in T_xM that is tangent to γ at x , pointing away from y , and that has length $d(x, y)$. This vector is $-\log_x y$, and hence (4.4) holds.

We begin with a result on time-dependent vector fields.

Lemma 4.1 (Fundamental Lemma I). *Let U be a totally normal neighbourhood of M and let X, Y be two time-dependent vector fields on U . Let $\Sigma \subset U$ be measurable and suppose that Ψ_X and Ψ_Y are flow maps defined on $\Sigma \times [0, \tau)$, for some $\tau > 0$, generated by (X, Σ) and (Y, Σ) , respectively. Assume furthermore that X is Lipschitz continuous with respect to its first variable, uniformly with respect to its second variable, i.e., there exists $L > 0$ such that*

$$\|X_t(x) - \Pi_{yx}X_t(y)\|_x \leq Ld(x, y), \quad \text{for all } (x, y, t) \in U \times U \times [0, \tau).$$

Then, for all $x \in \Sigma$, there holds

$$d(\Psi_X^t(x), \Psi_Y^t(x)) \leq \frac{e^{Lt} - 1}{L} \|X - Y\|_{L^\infty(U \times [0, \tau))}, \quad \text{for all } t \in [0, \tau).$$

Proof. Let $x \in \Sigma$ and, for a better readability of the following calculations, write $p = \Psi_X^t(x)$ and $q = \Psi_Y^t(x)$; note that $p, q \in U$. We can assume $p \neq q$ or the result is trivial.

Since U is totally normal, the function $t \mapsto d^2(p, q) = d^2(\Phi_X^t(x), \Phi_Y^t(x))$ is differentiable and, for all $t \in [0, \tau)$, we have

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} d^2(p, q) &= \frac{1}{2} \langle \nabla d_q^2(p), X_t(p) \rangle_p + \frac{1}{2} \langle \nabla d_p^2(q), Y_t(q) \rangle_q \\ &= - \langle \log_p(q), X_t(p) \rangle_p - \langle \log_q(p), Y_t(q) \rangle_q \\ &= - \langle \log_p(q), X_t(p) \rangle_p - \langle \Pi_{qp} \log_q(p), \Pi_{qp} Y_t(q) \rangle_p =: I + II, \end{aligned} \tag{4.5}$$

where for the third equality we used that the parallel transport is an isometry.

We add and subtract the quantities

$$A := \langle \Pi_{qp} \log_q(p), X_t(p) \rangle_p \quad \text{and} \quad B := \langle \Pi_{qp} \log_q(p), \Pi_{qp} X_t(q) \rangle_p$$

to the right-hand side of (4.5), which now reads

$$I - A + A - B + B + II.$$

The term $II + B$ estimates as

$$\begin{aligned} II + B &= \langle \Pi_{qp} \log_q(p), \Pi_{qp} X_t(q) - \Pi_{qp} Y_t(q) \rangle_p \\ &\leq \| \Pi_{qp} \log_q(p) \|_p \| \Pi_{qp} (X_t - Y_t)(q) \|_p \\ &= \| \log_q(p) \|_q \| X_t(q) - Y_t(q) \|_q \\ &\leq d(p, q) \| X - Y \|_{L^\infty(U \times [0, \tau])}, \end{aligned}$$

where we used again that Π_{qp} is an isometry and that $\| \log_q(p) \|_q = d(p, q)$. Also estimate

$$\begin{aligned} A - B &= \langle \Pi_{qp} \log_q(p), X_t(p) - \Pi_{qp} X_t(q) \rangle_p \\ &\leq \| \Pi_{qp} \log_q(p) \|_p \| X_t(p) - \Pi_{qp} X_t(q) \|_p \\ &\leq \| \log_q(p) \|_q L d(p, q) \\ &= L d(p, q)^2, \end{aligned}$$

where we used the Lipschitz continuity of X , the isometric property of Π_{qp} and that $\| \log_q(p) \|_q = d(p, q)$. Finally, by (4.4), we yield

$$I - A = - \langle \log_p(q) + \Pi_{qp} \log_q(p), X_t(p) \rangle_p = 0.$$

Using these estimates in (4.5), we find

$$\frac{1}{2} \frac{d}{dt} d^2(p, q) \leq L d^2(p, q) + \| X - Y \|_{L^\infty(U \times [0, \tau])} d(p, q),$$

and after cancelling a $d(p, q)$, we get:

$$\frac{d}{dt} d(p, q) \leq L d(p, q) + \| X - Y \|_{L^\infty(U \times [0, \tau])}.$$

Gronwall's lemma now gives the desired result. □

We now focus on what restrictions we should further impose on a neighbourhood U for the well-posedness to hold on U . Given a point $z \in M$, [29, Theorem 6.6.1] gives a bound on the Hessian of d_z^2 provided the sectional curvature \mathcal{K} of M is locally bounded. Specifically, let $r < \text{inj}(M)$ so that the exponential map \exp_z is a diffeomorphism on $B_r(0) \subset T_z M$. Denote by $B_r(z) := \exp_z(B_r(0))$ the geodesic ball

centred at z with radius r , and suppose that there exist reals $\lambda \leq 0$ and $0 \leq \mu \leq \frac{\pi^2}{4r^2}$ such that

$$\lambda \leq \mathcal{K} \leq \mu, \quad \text{on } B_r(z). \tag{4.6}$$

Then, by [29, Theorem 6.6.1] it holds that

$$2\sqrt{\mu} d(x, z) \cot(\sqrt{\mu} d(x, z)) \|v\|_x^2 \leq \langle \text{Hess}_v d_z^2(x), v \rangle_x \leq 2\sqrt{-\lambda} d(x, z) \coth(\sqrt{-\lambda} d(x, z)) \|v\|_x^2, \tag{4.7}$$

for all $x \in B_r(z)$ and $v \in T_x M$. Note that if M is simply connected, in addition to satisfying **(M)**, then $\text{inj}(M) = \infty$ whenever $\mathcal{K} \leq 0$ everywhere (cf. [29, Corollary 6.9.1], a consequence of the Cartan–Hadamard theorem).

Let us further introduce some notation used throughout the rest of the paper:

Notation. Because M is locally compact (since finite-dimensional) and \mathcal{K} is continuous, \mathcal{K} is also locally bounded (cf. [33, Remark 2.2]). For any $p \in M$ and $r > 0$, we write $\lambda_r(p) \leq 0$ and $\mu_r(p) \geq 0$ any lower and upper bounds for \mathcal{K} on the set $\{x \in M \mid d(p, x) < r\}$ (which, by the terminology set up in Appendix A, is denoted by $B_r(p)$ if $r \leq r_{\text{inj}}(p)$).

Lemma 4.2 (Fundamental Lemma II). *Let $U \subset B_{r/2}(p)$ be open for some $p \in M$ and $r < \text{inj}(M)$, and denote $\lambda = \lambda_{2r}(p)$ and $\mu = \mu_{2r}(p)$. Suppose moreover that U satisfies*

$$\Delta := \text{diam}(U) < \frac{\pi}{2\sqrt{\mu}}, \quad \text{if } \mu > 0. \tag{4.8}$$

Then, for all $z \in U$, the Hessian $\text{Hess}_z d_z^2$ is bounded on U by $L := 2\sqrt{-\lambda} \Delta \coth(\sqrt{-\lambda} \Delta)$.

If moreover U is geodesically convex, we have

$$\|\log_x z - \Pi_{yx} \log_y z\|_x \leq \frac{L}{2} d(x, y), \quad \text{for all } x, y, z \in U. \tag{4.9}$$

Proof. The proof is direct from (4.7); see Appendix B.5 for the details. □

Remark 1. When $\lambda_r(p) = \mu_r(p) = 0$ for all $p \in M$ and $r > 0$, we recover the trivial result that $L = 2$ in Lemma 4.2 (and, in particular, L is independent of the diameter of the subset U), and (4.9) holds for all $x, y, z \in M$.

Lemma 4.2 states that, for all z , the map $x \mapsto \log_x z$ is Lipschitz continuous in the sense of parallel transport. On the other hand, the following lemma shows that, for all x , the map $z \mapsto \log_x z$ is Lipschitz continuous in the classical Euclidean sense.

Lemma 4.3 (Fundamental Lemma III). *Let $U \subset B_r(p)$ be a totally normal neighbourhood for some $p \in M$ and $r < \text{inj}(M)$, and write $\mu = \mu_r(p)$. We distinguish two cases:*

(i) $\mu = 0$. *Then,*

$$\|\log_x y - \log_x z\|_x \leq d(y, z), \quad \text{for all } x, y, z \in U. \tag{4.10}$$

(ii) $\mu > 0$. *Take any $0 < \varepsilon < \pi$ and further assume $\text{diam}(U) \leq \frac{\pi - \varepsilon}{\sqrt{\mu}}$. Then,*

$$\|\log_x z - \log_x y\|_x \leq \frac{\pi - \varepsilon}{\sin(\pi - \varepsilon)} d(y, z), \quad \text{for all } x, y, z \in U. \tag{4.11}$$

Proof. This follows from [29, Corollary 6.6.1]; see the details in Appendix B.6. □

Remark 2. When $\text{inj}(M) = \infty$ and $\mu_r(p) = 0$ for all $p \in M$ and $r > 0$, i.e., M is globally nonpositively curved, (4.10) holds for all $x, y, z \in M$.

Remark 3. If M is globally nonpositively curved, then the restriction on U in Lemma 4.2 in order to get a bounded Hessian simplifies into open and bounded. In Lemma 4.3 the restriction on U simplifies into totally normal neighborhood and bounded.

4.2. Local well-posedness. Let us now turn to the local well-posedness of solutions to Equation (1.1). We consider an additional assumption:

(Kloc) $K: M \times M \rightarrow \mathbb{R}$ satisfies **(K)** with g' locally Lipschitz continuous.

For $\Delta > 0$, we write $C_{g'}(\Delta)$ and $L_{g'}(\Delta)$ for the L^∞ norm and the Lipschitz constant of g' on $[0, \Delta^2]$.

Lemma 4.4. Assume the manifold M and potential K satisfy **(M)** and **(Kloc)**, and take $U \subset B_{r/2}(p)$ open and geodesically convex for some $p \in M$ and $r < \text{inj}(M)$ such that (4.8) holds (cf. also Notation 4.1). Let $T \in (0, \infty]$ and $\rho \in C([0, T]; \mathcal{P}(U))$. Then the time-dependent vector field $v[\rho]$ given by (3.1) is bounded and Lipschitz continuous with respect to its first variable, uniformly with respect to its second variable.

Proof. Write Δ for $\text{diam}(U)$. We have, for all $x, y \in U$,

$$\|\nabla K_y(x)\|_x = \|g'(d(x, y)^2)\nabla d_y^2(x)\|_x \leq 2C_{g'}(\Delta)\Delta,$$

where we used the local bound on g' and that $\|\nabla d_y^2(x)\|_x = 2d(x, y) \leq 2\Delta$. Immediately, for all $t \in [0, T]$, it follows

$$\|v[\rho](x, t)\|_x \leq \int_U \|\nabla K_y(x)\|_x \, d\rho_t(y) \leq 2C_{g'}(\Delta)\Delta, \tag{4.12}$$

which shows that $v[\rho]$ is bounded.

Also, for all $x, y, z \in U$, we get

$$\begin{aligned} & \|\nabla K_z(x) - \Pi_{yx}\nabla K_z(y)\|_x \\ &= \|g'(d(x, z)^2)\nabla d_z^2(x) - g'(d(y, z)^2)\Pi_{yx}\nabla d_z^2(y)\|_x \\ &\leq |g'(d(x, z)^2)| \|\nabla d_z^2(x) - \Pi_{yx}\nabla d_z^2(y)\|_x + \|\nabla d_z^2(y)\|_y |g'(d(x, z)^2) - g'(d(y, z)^2)| \\ &\leq C_{g'}(\Delta)L d(x, y) + 2d(y, z)L_{g'}(\Delta)|d(x, z)^2 - d(y, z)^2|, \end{aligned}$$

where we used the Lipschitz continuity of the vector field ∇d_z^2 (equivalently of the logarithm map) given by (4.9) and the local bound and Lipschitz continuity of g' . Now use $d(y, z) \leq \Delta$ and

$$|d(x, z)^2 - d(y, z)^2| = |d(x, z) - d(y, z)||d(x, z) + d(y, z)| \leq 2\Delta d(x, y),$$

to get

$$\|\nabla K_z(x) - \Pi_{yx}\nabla K_z(y)\|_x \leq \bar{L} d(x, y) \quad \text{for all } x, y, z \in U,$$

where

$$\bar{L} = C_{g'}(\Delta)L + 4\Delta^2 L_{g'}(\Delta). \tag{4.13}$$

Then, for all $t \in [0, T]$ and $x, y \in U$, we get

$$\begin{aligned} \|v[\rho](x, t) - \Pi_{yx}v[\rho](y, t)\|_x &\leq \int_U \|\nabla K_z(x) - \Pi_{yx}\nabla K_z(y)\|_x \, d\rho_t(z) \\ &\leq \bar{L} d(x, y), \end{aligned} \tag{4.14}$$

which shows that $v[\rho]$ is Lipschitz continuous with respect to its first variable (space), uniformly with respect to its second (time). \square

Remark 4. The L^∞ bound and the Lipschitz constant of $v[\rho]$ in (4.12) and (4.14) do not depend on the curve ρ . This is important for the proof of Theorem 4.6.

Lemma 4.5. *Let M and K satisfy **(M)** and **(Kloc)** and $U \subset B_r(p)$ be a totally normal neighborhood for some $p \in M$ and $r < \text{inj}(M)$, and write $\mu = \mu_r(p)$ (recall Notation 4.1). If $\mu > 0$ assume that $\Delta := \text{diam}(U)$ satisfies $\Delta \leq \frac{\pi-\varepsilon}{\sqrt{\mu}}$ for some $0 < \varepsilon < \pi$. Let $\rho, \sigma \in C([0, T]; \mathcal{P}(U))$. Then, there exists $\Lambda > 0$ so that*

$$\|v[\rho] - v[\sigma]\|_{L^\infty(U \times [0, T])} \leq \Lambda \mathbf{W}_1(\rho, \sigma).$$

Proof. By Lemma 4.3, we have

$$\|\log_x z - \log_x y\|_x \leq \ell d(y, z) \quad \text{for all } x, y, z \in U,$$

where $\ell = \frac{\pi-\varepsilon}{\sin(\pi-\varepsilon)}$ if $\mu > 0$ and $\ell = 1$ if $\mu = 0$. Hence, using (4.2) we find

$$\|\nabla d_y^2(x) - \nabla d_z^2(x)\|_x \leq 2\ell d(y, z), \quad \text{for all } x, y, z \in U.$$

Since **(Kloc)** holds, for all $x, y \in U$ compute

$$\begin{aligned} & \|\nabla K_y(x) - \nabla K_z(x)\|_x \\ &= \|g'(d(x, y)^2) \nabla d_y^2(x) - g'(d(x, z)^2) \nabla d_z^2(x)\|_x \\ &\leq |g'(d(x, z)^2)| \|\nabla d_y^2(x) - \nabla d_z^2(x)\|_x + \|\nabla d_y^2(x)\|_x |g'(d(x, y)^2) - g'(d(x, z)^2)|, \end{aligned}$$

where we added and subtracted $g'(d(x, z)^2) \nabla d_y^2(x)$ on the first line and used the triangle inequality. Then, using the bound and Lipschitz constant of g' , the fact that $\|\nabla d_y^2(x)\|_x = 2d(x, y)$ and again the triangle inequality, we find

$$\begin{aligned} & \|\nabla K_y(x) - \nabla K_z(x)\|_x \\ &\leq 2C_{g'}(\Delta) \ell d(y, z) + 2L_{g'}(\Delta) |d(x, y) + d(x, z)| |d(x, y) - d(x, z)| d(x, y) \\ &\leq (2C_{g'}(\Delta) \ell + 4L_{g'}(\Delta) \Delta^2) d(y, z). \end{aligned} \tag{4.15}$$

Now, for $(x, t) \in U \times [0, T]$, take $\pi_t \in \Pi(\rho_t, \sigma_t)$ to be an optimal transport plan between ρ_t and σ_t , and estimate

$$\begin{aligned} \|v[\rho](x, t) - v[\sigma](x, t)\|_x &= \left\| \int_U \nabla K_y(x) \, d\rho_t(y) - \int_U \nabla K_z(x) \, d\sigma_t(z) \right\|_x \\ &= \left\| \int_{U \times U} \nabla K_y(x) \, d\pi_t(y, z) - \int_{U \times U} \nabla K_z(x) \, d\pi_t(y, z) \right\|_x \\ &\leq \int_{U \times U} \|\nabla K_y(x) - \nabla K_z(x)\|_x \, d\pi_t(y, z). \end{aligned}$$

Then, using (4.15) we find

$$\begin{aligned} \|v[\rho](x, t) - v[\sigma](x, t)\|_x &\leq \Lambda \int_{U \times U} d(y, z) \, d\pi_t(y, z) = \Lambda W_1(\rho_t, \sigma_t) \\ &\leq \Lambda \mathbf{W}_1(\rho, \sigma), \end{aligned} \tag{4.16}$$

where

$$\Lambda = 2C_{g'}(\Delta) \ell + 4L_{g'}(\Delta) \Delta^2. \tag{4.17}$$

Taking the supremum in $(x, t) \in U \times [0, T]$ on the left-hand side of (4.16) gives the result. \square

Remark 5. The assumptions on U in Lemma 4.5 are weaker than in Lemma 4.4.

Theorem 4.6 (Local Well-Posedness). *Assume the manifold M and interaction potential K satisfy **(M)** and **(Kloc)**, and take $U \subset B_{r/2}(p)$ open and geodesically convex for some $p \in M$ and $r < \text{inj}(M)$ so that (4.8) holds (cf. Notation 4.1). Let $\rho_0 \in \mathcal{P}(U)$. Then, there exist $T > 0$ and a unique weak solution in $C([0, T]; \mathcal{P}(U))$ starting from ρ_0 to the aggregation equation (1.1).*

Proof. By Lemma 4.4 the velocity field $v[\sigma]$ (for any fixed continuous curve σ in $\mathcal{P}(U)$) satisfies the assumptions of Theorem 2.5, so that we can infer that $(v[\sigma], \text{supp}(\rho_0))$ generates a unique maximal flow $\Psi_{v[\sigma]}$ defined on $\text{supp}(\rho_0) \times [0, \tau)$ for some $\tau \in (0, \infty]$. In addition, since the L^∞ bound and the Lipschitz constant in (4.12) and (4.14) do not depend on the underlying curve (see Remark 4), the maximal time of existence τ of $\Psi_{v[\sigma]}$ does not depend on σ by the Cauchy–Lipschitz theorem (cf. Theorem 2.5).

We can then define the map Γ by

$$\Gamma(\sigma)(t) = \Psi_{v[\sigma]}^t \# \rho_0, \quad \text{for all } \sigma \in C([0, \tau]; \mathcal{P}(U)) \text{ and } t \in [0, \tau).$$

We will prove that Γ defines a map from $C([0, \tau]; \mathcal{P}(U))$ into itself and that it has a unique fixed point. This fixed point is the weak solution of (1.1) starting at ρ_0 .

Note first that Γ does indeed map $C([0, \tau]; \mathcal{P}(U))$ into itself. For $\sigma \in C([0, \tau]; \mathcal{P}(U))$ fixed, by definition of a flow map generated by $(v[\sigma], \text{supp}(\rho_0))$, we infer that $\Psi_{v[\sigma]}^t(x) \in U$ for all $x \in \text{supp}(\rho_0)$ and $t \in [0, \tau)$. Consequently, $\Gamma(\sigma)(t)$ is supported in U and moreover, it is a probability measure by conservation of mass through the push-forward, so that $\Gamma(\sigma)(t) \in \mathcal{P}(U)$, for all $t \in [0, \tau)$. Furthermore, the map $t \mapsto \Gamma(\sigma)(t)$ is continuous due to the combination of Lemma 3.2(ii) with Lemma 4.4. We conclude that $\Gamma: (C([0, \tau]; \mathcal{P}(U)), \mathbf{W}_1) \rightarrow (C([0, \tau]; \mathcal{P}(U)), \mathbf{W}_1)$.

To show that Γ is a contraction we will have to restrict the final time to some $T \leq \tau$ to be determined. Let $\rho, \sigma \in C([0, \tau]; \mathcal{P}(U))$. Then, for all $t \in [0, \tau)$, we have

$$\begin{aligned} W_1(\Psi_{v[\rho]}^t \# \rho_0, \Psi_{v[\sigma]}^t \# \rho_0) &\leq \sup_{x \in \text{supp}(\rho_0)} d(\Psi_{v[\rho]}^t(x), \Psi_{v[\sigma]}^t(x)) \\ &\leq C(t) \|v[\rho] - v[\sigma]\|_{L^\infty(U \times [0, \tau])} \\ &\leq C(t) \Lambda W_1(\rho, \sigma), \end{aligned} \tag{4.18}$$

where for the first inequality we used Lemma 3.2(i), for the second inequality we used Lemmas 4.4 and 4.1, with

$$C(t) = \frac{e^{\bar{L}t} - 1}{\bar{L}},$$

and for the last inequality we used Lemma 4.5; the Lipschitz constants \bar{L} and Λ depend on $\text{diam}(U)$ and are defined in (4.13) and (4.17). Since $t \mapsto C(t)$ is increasing, with $\lim_{t \rightarrow 0} C(t) = 0$ and Λ is independent of time, we can choose $T \leq \tau$ (only depending on \bar{L} and Λ) small enough so that

$$C(t)\Lambda < C(T)\Lambda < \bar{C}, \quad \text{for all } t \in [0, T),$$

for some constant $\bar{C} < 1$. Restricting T accordingly, by taking the supremum over $[0, T)$ in (4.18) we find

$$W_1(\Gamma(\rho), \Gamma(\sigma)) \leq \bar{C} W_1(\rho, \sigma),$$

with $\bar{C} < 1$. This shows that the restriction of Γ to $(C([0, T]; \mathcal{P}(U)), \mathbf{W}_1)$ is a contraction.

We have thus shown that $\Gamma: (C([0, T]; \mathcal{P}(U)), \mathbf{W}_1) \rightarrow (C([0, T]; \mathcal{P}(U)), \mathbf{W}_1)$ has a unique fixed point, that is, there exists a unique $\rho \in C([0, T]; \mathcal{P}(U))$ such that

$$\rho_t = \Psi_{v[\rho]}^t \# \rho_0 \quad \text{for all } t \in [0, T];$$

this fixed point ρ is the desired solution. □

Remark 6. In Theorem 4.6, one could choose $U = B_{r/2}(p)$ as long as this geodesic ball is convex and $r < \text{inj}(M)$, that is, as long as $r < \min(\text{inj}(M), 2 \text{conv}(M)) = 2 \text{conv}(M)$ (where the equality comes from the fact that $2 \text{conv}(M) \leq \text{inj}(M)$ [11, Proposition IX.6.1]).

4.3. Global well-posedness when M is nonpositively curved. We establish here the global well-posedness for (1.1) when M is nonpositively curved. Recall that in this case, if M is simply connected, then $\text{inj}(M) = \infty$ and M is unbounded; in fact, we have that the exponential map is a diffeomorphism from $T_x M$ to M for all $x \in M$ (see [29, Corollary 6.9.1]), so that, in particular, the velocity field (3.1) is well-defined for all $\rho \in C([0, \infty); \mathcal{P}_\infty(M))$ and M is geodesically convex.

We shall focus on the case when M is nonpositively curved with bounded curvature. We thus consider the following hypothesis, for some $\lambda \leq 0$:

(M) $_\lambda$ M satisfies **(M)**, is simply connected, and its curvature \mathcal{K} is such that $\lambda \leq \mathcal{K} \leq 0$ everywhere.

We wish to adapt Lemmas 4.4 and 4.5 globally when M satisfies **(M) $_\lambda$** for some $\lambda \leq 0$. To this end, consider the following hypothesis on K , for some $\lambda \leq 0$:

(Kglob) $_\lambda$ $K: M \times M \rightarrow \mathbb{R}$ satisfies **(K)** with g so that, for some constant $A_{g'} > 0$,

$$|g'(r^2)r - g'(s^2)s| \leq A_{g'} |r - s|, \quad \text{for all } r, s \geq 0. \tag{4.19}$$

When $\lambda < 0$, it further holds that

$$|g'(r^2)|r \leq A_{g'}, \quad \text{for all } r \geq 0.$$

The first bound in **(Kglob) $_\lambda$** means that the function $r \mapsto g'(r^2)r$ is globally Lipschitz continuous; it implies in particular that g' is globally bounded by $A_{g'}$. The second bound in **(Kglob) $_\lambda$** , in case $\lambda < 0$, means that $r \mapsto g'(r^2)r$ is globally bounded; it also says that $g'(r)$ must decrease at least as fast as $1/\sqrt{r}$ as $r \rightarrow \infty$.

The Lipschitz condition in **(Kglob) $_\lambda$** is consistent with the classical Cauchy–Lipschitz theory in Euclidean space. Indeed, in this case this Lipschitz condition is enough to get the global Lipschitz continuity of the velocity field (and the local boundedness that follows from it) in order to obtain global well-posedness of the flow maps (cf. Theorem 2.6). When the space is negatively curved, the Lipschitz condition in **(Kglob) $_\lambda$** is not sufficient anymore to get the global Lipschitz continuity of the velocity field. This stems from (4.9), where the Lipschitz constant depends on the distance between the points considered (as opposed to the flat case discussed in Remark 1)—the global upper bound condition in **(Kglob) $_\lambda$** helps counterbalance this effect. All this can be seen in details in the proof of Lemma 4.7 below.

Remark 7. One can check that the Lipschitz condition in **(Kglob) $_\lambda$** implies

$$|g'(r^2) - g'(s^2)|s \leq 2A_{g'} |r - s|, \quad \text{for all } r, s \geq 0.$$

Lemma 4.7. Assume that there exists $\lambda \leq 0$ such that M satisfies **(M) $_\lambda$** and K satisfies **(Kglob) $_\lambda$** . Let $\rho \in C([0, \infty); \mathcal{P}_\infty(M))$. Then the time-dependent vector

field $v[\rho]$ given by (3.1) is locally bounded and globally Lipschitz continuous with respect to its first variable, uniformly with respect to its second variable.

Proof. We follow closely the proof of Lemma 4.4. Let $U \subset M$ be bounded, open and convex with $\Delta = \text{diam}(U)$. We have, for all $x, y \in U$,

$$\|\nabla K_y(x)\|_x = \|g'(d(x, y)^2)\nabla d_y^2(x)\|_x \leq 2A_{g'}\Delta,$$

which yields the local boundedness of $v[\rho](\cdot, t)$, uniformly with respect to $t \in [0, \infty)$. (Note that in case $\lambda < 0$, this bound on $v[\rho](\cdot, t)$ is actually global thanks to the second condition in **(Kglob)** $_\lambda$.)

For all $x, y, z \in U$, we get

$$\begin{aligned} & \|\nabla K_z(x) - \Pi_{yx}\nabla K_z(y)\|_x \\ &= \|g'(d(x, z)^2)\nabla d_z^2(x) - g'(d(y, z)^2)\Pi_{yx}\nabla d_z^2(y)\|_x \\ &\leq |g'(d(x, z)^2)|\|\nabla d_z^2(x) - \Pi_{yx}\nabla d_z^2(y)\|_x + \|\nabla d_z^2(y)\|_y |g'(d(x, z)^2) - g'(d(y, z)^2)| \\ &\leq |g'(\Delta^2)|L(\Delta)d(x, y) + 4A_{g'}|d(x, z) - d(y, z)| \\ &\leq (|g'(\Delta^2)|L(\Delta) + 4A_{g'})d(x, y), \end{aligned}$$

where for the first inequality we added and subtracted $g'(d(x, z)^2)\Pi_{yx}\nabla d_z^2(y)$ and used the triangle inequality; we then used the Lipschitz continuity of the vector field ∇d_z^2 given by (4.9) with $L(\Delta) := 2\sqrt{-\lambda}\Delta \coth(\sqrt{-\lambda}\Delta)$, the fact that $\|\nabla d_z^2(y)\|_y = 2d(y, z)$ and Remark 7 for the second inequality, and the reverse triangle inequality on d for the third inequality. We notice that if $\lambda = 0$, then $L(\Delta) = 2$ and so $|g'(\Delta^2)|L(\Delta) \leq 2A_{g'}$. If $\lambda < 0$, we have

$$\begin{aligned} g'(\Delta^2)L(\Delta) &= 2g'(\Delta^2)\sqrt{-\lambda}\Delta \coth(\sqrt{-\lambda}\Delta) \\ &\leq \begin{cases} 2g'(\Delta^2) \coth(1) & \text{if } \Delta \leq \frac{1}{\sqrt{-\lambda}} \\ 2g'(\Delta^2)\sqrt{-\lambda}\Delta \coth(1) & \text{if } \Delta > \frac{1}{\sqrt{-\lambda}} \end{cases} \\ &\leq \begin{cases} 2A_{g'} \coth(1) & \text{if } \Delta \leq \frac{1}{\sqrt{-\lambda}} \\ 2A_{g'}\sqrt{-\lambda} \coth(1) & \text{if } \Delta > \frac{1}{\sqrt{-\lambda}} \end{cases} \\ &\leq 2A_{g'} \coth(1) \max(1, \sqrt{-\lambda}), \end{aligned}$$

where for the first inequality we used that $r \coth(r)$ and $\coth(r)$ are bounded by $\coth(1)$, respectively for $r \leq 1$ and for $r > 1$, and for the second inequality we used the second bound condition in **(Kglob)** $_\lambda$. Thus, because $\coth(1) \max(1, \sqrt{-\lambda}) \geq 1$, in either case $\lambda = 0$ or $\lambda < 0$ we get

$$g'(\Delta^2)L(\Delta) \leq 2A_{g'} \coth(1) \max(1, \sqrt{-\lambda}).$$

Since the right-hand side above is independent of U , and U is any bounded, open, convex subset of M , we deduce that for all $x, y, z \in M$ we have

$$\|\nabla K_z(x) - \Pi_{yx}\nabla K_z(y)\|_x \leq 2A_{g'}(\coth(1) \max(1, \sqrt{-\lambda}) + 2)d(x, y),$$

We then conclude with Lipschitz constant $\bar{L} = 2A_{g'}(\coth(1) \max(1, \sqrt{-\lambda}) + 2)$. \square

Lemma 4.8. *Suppose that M satisfies **(M)** $_\lambda$ for some $\lambda \leq 0$ and K satisfies the Lipschitz condition (4.19). Let furthermore $\rho, \sigma \in C([0, \infty); \mathcal{P}_\infty(M))$. Then, there*

exists $\Lambda > 0$ so that

$$\|v[\rho] - v[\sigma]\|_{L^\infty(M \times [0, \infty))} \leq \Lambda \mathbf{W}_1(\rho, \sigma).$$

Proof. We follow closely the proof of Lemma 4.5. By Lemma 4.3 and Remark 2, recall that

$$\|\nabla d_y^2(x) - \nabla d_z^2(x)\|_x \leq 2d(y, z), \quad \text{for all } x, y, z \in M. \quad (4.20)$$

For all $x, y, z \in M$ compute

$$\begin{aligned} & \|\nabla K_y(x) - \nabla K_z(x)\|_x \\ &= \|g'(d(x, y)^2) \nabla d_y^2(x) - g'(d(x, z)^2) \nabla d_z^2(x)\|_x \\ &\leq |g'(d(x, z)^2)| \|\nabla d_y^2(x) - \nabla d_z^2(x)\|_x + \|\nabla d_y^2(x)\|_x |g'(d(x, y)^2) - g'(d(x, z)^2)| \\ &\leq 2A_{g'} d(y, z) + 4A_{g'} |d(x, y) - d(x, z)| \\ &\leq 6A_{g'} d(y, z), \end{aligned}$$

where for the first inequality we added and subtracted $g'(d(x, y)^2) \nabla d_z^2(x)$ and used triangle inequality; we then used (4.20), the fact that $\|\nabla d_y^2(x)\|_x = 2d(x, y)$ and Remark 7 for the second inequality, and the reverse triangle inequality on d for the third inequality. Following the proof of Lemma 4.5, we conclude with $\Lambda = 6A_{g'}$. \square

In addition to global well-posedness, the following theorem shows that the support of the global solution is contained in an increasing geodesic ball centred around the initial data.

Theorem 4.9 (Global Well-Posedness on Nonpositively Curved Manifold). *Assume that there exists $\lambda \leq 0$ such that M satisfies $(\mathbf{M})_\lambda$ and K satisfies $(\mathbf{Kglob})_\lambda$. Let $\rho_0 \in \mathcal{P}_\infty(M)$. Then, there exists a unique weak solution in $C([0, \infty); \mathcal{P}_\infty(M))$ starting from ρ_0 to the aggregation equation (1.1). Moreover, there exist a nondecreasing function $R: [0, \infty) \rightarrow \mathbb{R}$ and a point $p \in \text{supp}(\rho_0)$ such that*

$$\text{supp}(\rho_t) \subset B_{R(t)}(p), \quad \text{for all } t \geq 0.$$

Proof. By Lemma 4.7, since M is convex, the velocity field $v[\sigma]$ (for any fixed continuous curve σ in $\mathcal{P}_\infty(M)$) satisfies the assumptions of the global Cauchy-Lipschitz theorem (cf. Theorem 2.6), so that we can infer that $(v[\sigma], \text{supp}(\rho_0))$ generates a unique global flow defined on $\text{supp}(\rho_0) \times [0, \infty)$.

Following the same steps as in the proof of Theorem 4.6, using Lemmas 4.7 and 4.8 in place of Lemmas 4.4 and 4.5, respectively, yields the well-posedness by observing that the constants \bar{L} and Λ are now independent of $\text{supp}(\rho_0)$ and its diameter, which allows us to apply an iterative argument indefinitely and obtain a unique solution in $C([0, \infty); \mathcal{P}_\infty(M))$. The boundedness of the support of the solution and the existence of a nondecreasing ball containing it follows from the iterative construction of the global flow map in the proof of Theorem 2.6. \square

5. Additional results. We give here two additional results which can be derived by following very closely the global well-posedness and stability theory established for the model on the sphere in [20]. For this reason, we do not give the proofs here and refer the reader to [20, Proposition 4.1] (global well-posedness for attractive potentials) and [20, Theorem 3.8] (stability of solutions) should they wish to see details.

Following [20], we give the following definition of purely attractive potential: given $x \in M$, we say that K is *purely attractive* at x if for all normal neighbourhood U of x and $y \in U$ we have

$$\langle \nabla K_y(x), \log_x y \rangle_x \leq 0.$$

Consider now the hypothesis below on K :

(Katt) $K : M \times M \rightarrow \mathbb{R}$ satisfies **(Kloc)** with g such that

$$g'(r^2) \geq 0 \quad \text{for all } r < \text{inj}(M).$$

By (4.1), the nonnegativity condition in **(Katt)** means that K is purely attractive.

Theorem 5.1 (Global Well-Posedness for Attractive Potential). *Let M and K satisfy **(M)** and **(Katt)**, and let $U \subset B_{r/2}(p)$ for some $p \in M$ and $r < \text{inj}(M)$ be open and geodesically convex and such that (4.8) holds. Let moreover $\rho_0 \in \mathcal{P}(U)$ and $R > 0$ be such that $\text{supp}(\rho_0) \subset B_R(p) \subset U$. Then, there exists a unique weak solution ρ in $C([0, \infty); \mathcal{P}(U))$ starting from ρ_0 to the aggregation equation (1.1); furthermore, $\text{supp}(\rho_t) \subset \overline{B_R(p)}$ for all $t \geq 0$.*

Theorem 5.2 (Stability). *Let $\rho_0, \sigma_0 \in \mathcal{P}_\infty(M)$. Let furthermore ρ and σ be weak solutions to (1.1) defined on some time interval $[0, T)$ starting from ρ_0 and σ_0 , respectively, as derived in Theorems 4.6, 4.9 or 5.1. Then,*

$$W_1(\rho_t, \sigma_t) \leq e^{(\bar{L} + \Lambda)t} W_1(\rho_0, \sigma_0) \quad \text{for all } t \in [0, T),$$

where \bar{L} and Λ are the constants given in the proofs of Lemmas 4.4 and 4.5 in the local, compact and purely attractive cases and of Lemmas 4.7 and 4.8 in the global, unbounded case.

We note that when M and K verify **(M) $_\lambda$** and **(Kglob) $_\lambda$** for some $\lambda \leq 0$, the above theorem illustrates the fact that the larger (in magnitude) the lower bound λ on \mathcal{K} is, the faster solutions may spread apart in time.

Appendix A. Basic concepts and terminology from differential geometry.

All the concepts discussed here are standard and can be found in any graduate differential geometry book. We refer the reader for example to [11, 13, 31].

Logarithm map and normal neighbourhoods. Given $x \in M$, there exists $G \subset T_x M$ open with $0 \in G$ such that the exponential map $\exp_x : G \rightarrow U := \exp_x(G)$ restricted to G is a diffeomorphism. The inverse of \exp_x is the *logarithm map* at x (on U), denoted by $\log_x : U \rightarrow G$. We have that $x \in U$ and U is an open subset of M . We call any such U a *normal neighbourhood* of x —conversely, any $x \in U$ has a normal neighbourhood. A subset of M is said to be a *totally normal neighbourhood* if it is a normal neighbourhood of each of its points.

If $r > 0$ is such that \exp_x defines a diffeomorphism on the open ball $B_r(0) \subset T_x M$, then $B_r(x) := \exp_x(B_r(0))$ is open and is called the *open geodesic ball centred at x of radius r* ; by the previous paragraph, $B_r(x)$ is by definition a normal neighbourhood of x . Also, it coincides with the metric ball centred at x with radius r , that is, with the set $\{y \in M \mid d(x, y) < r\}$. The largest such radius r is called the *injectivity radius at x* and is denoted $r_{\text{inj}}(x)$. The infimum over all $y \in M$ of $r_{\text{inj}}(y)$ is referred to as the *injectivity radius of M* , which we denote by $\text{inj}(M)$.

Geodesics and convexity. By definition of the exponential map, any point y in a normal neighbourhood of another point x can be connected to x via a unique

minimizing geodesic $\gamma_{xy}: [0, 1] \rightarrow M$, which satisfies, for all $t \in [0, 1]$,

$$\gamma_{xy}(t) = \exp_x(t \log_x y).$$

This formula indicates that \exp_x transforms the straight line $t \mapsto t \log_x(y)$ in $T_x M$ into the geodesic connecting x and y .

A (*geodesically*) *convex* subset of M is a subset in which each pair of points can be connected by a unique minimizing geodesic which lies entirely within it—when open, it is thus a particular case of a totally normal neighbourhood that minimizing geodesics cannot exit. We say that M is (*geodesically*) *convex* if it is a geodesically convex subset of itself. Note that a geodesic ball is not necessarily convex. Nevertheless, given any $x \in M$, it is a fact that x has a convex geodesic ball as a neighbourhood; we refer to the radius of the largest such ball as the *convexity radius at x* , denoted $r_{\text{conv}}(x)$. The infimum over all $y \in M$ of $r_{\text{conv}}(y)$ is referred to as the *convexity radius of M* , which we denote by $\text{conv}(M)$.

Normal charts. Let $x \in M$ and U be a normal neighbourhood of x . We can construct a chart in a canonical way. Indeed, let $\{E_1(x), \dots, E_n(x)\}$ be an orthonormal basis of $T_x M$ and define $E: T_x M \rightarrow \mathbb{R}^n$ by

$$E(v) = (v_1, \dots, v_n), \quad \text{for all } v = v_1 E_1(x) + \dots + v_n E_n(x) \in T_x M, \quad (\text{A.1})$$

i.e., $E(v)$ gives the vector in \mathbb{R}^n whose coordinates are the coordinates of v in $\{E_1(x), \dots, E_n(x)\}$. Then, define the map $\varphi: U \rightarrow \mathbb{R}^n$ by

$$\varphi(y) = E \circ \log_x(y) \quad \text{for all } y \in U;$$

note that $\varphi(x) = 0$. The pair (U, φ) induces a local chart of M containing x . Any chart thus constructed is referred to as a *normal chart* (generated by x).

Push-forward and coordinates. Given a differentiable function $f: M \rightarrow \mathbb{R}^n$, a point $x \in M$ and $v \in T_x M$, we call $f_* v := df(x)(v) \in \mathbb{R}^n$ the *push-forward* of v through f . Given any chart (V, φ) of M containing some $x, y \in M$, we write

$$\left\{ \frac{\partial}{\partial \varphi^1}(y), \dots, \frac{\partial}{\partial \varphi^n}(y) \right\}$$

for the basis of $T_y M$ defined by

$$\frac{\partial}{\partial \varphi^i}(y) = d\varphi^{-1}(\varphi(y))(e_i), \quad \text{that is, } e_i = d\varphi(y) \left(\frac{\partial}{\partial \varphi^i}(y) \right) = \varphi_* \frac{\partial}{\partial \varphi^i}(y),$$

where, for all $i \in \{1, \dots, n\}$, e_i is the i th vector of the canonical basis of \mathbb{R}^n . For all $v \in T_y M$ there exist $v_1, \dots, v_n \in \mathbb{R}$ such that

$$v = v_1 \frac{\partial}{\partial \varphi^1}(y) + \dots + v_n \frac{\partial}{\partial \varphi^n}(y).$$

Then, by linearity,

$$\begin{aligned} \varphi_* v &= d\varphi(y)(v) = v_1 d\varphi(y) \left(\frac{\partial}{\partial \varphi^1}(y) \right) + \dots + v_n d\varphi(y) \left(\frac{\partial}{\partial \varphi^n}(y) \right) \\ &= v_1 e_1 + \dots + v_n e_n = (v_1, \dots, v_n). \end{aligned}$$

Hence $\varphi_* v$ gives the vector in \mathbb{R}^n whose coordinates are those of v in the basis $\{\partial/\partial \varphi^1(y), \dots, \partial/\partial \varphi^n(y)\}$.

Euclidean norm and intrinsic distance. We give now useful relations between norms in \mathbb{R}^n and intrinsic distances on M . Let (V, φ) be a chart of M and let $Q \subset V$

be compact. Then there exist $c_{\varphi,Q}, C_{\varphi,Q} > 0$ so that

$$c_{\varphi,Q} \|\varphi_*v\|_{\mathbb{R}^n} \leq \|v\|_x \leq C_{\varphi,Q} \|\varphi_*v\|_{\mathbb{R}^n}, \quad \text{for all } x \in Q \text{ and } v \in T_xM. \quad (\text{A.2})$$

If furthermore $\varphi(U) \subset \mathbb{R}^n$ is convex, then there exists $L_{\varphi,Q} > 0$ so that

$$d(x, y) \leq L_{\varphi,Q} \|\varphi(x) - \varphi(y)\|_{\mathbb{R}^n}, \quad \text{for all } x, y \in Q. \quad (\text{A.3})$$

We now prove these relations. Let us first show (A.2). Define

$$T_\varphi = \{(x, v) \in Q \times TM \mid v \in T_xM, \|\varphi_*v\|_{\mathbb{R}^n} = 1\},$$

which is compact as well. Then, because $(x, v) \mapsto \|v\|_x$ is continuous on T_φ and we have $\|v\|_x > 0$ for all $(x, v) \in T_\varphi$, there exist $c_{\varphi,Q}, C_{\varphi,Q} > 0$ such that

$$c_{\varphi,Q} \leq \|v\|_x \leq C_{\varphi,Q}, \quad \text{for all } (x, v) \in T_\varphi.$$

Fix now $(x, v) \in Q \times TM$ with $v \in T_xM$ and $\|\varphi_*v\|_{\mathbb{R}^n} \neq 0$. Using the inequality above we get

$$\|v\|_x = \|\varphi_*v\|_{\mathbb{R}^n} \left\| \frac{v}{\|\varphi_*v\|_{\mathbb{R}^n}} \right\|_x \leq C_{\varphi,Q} \|\varphi_*v\|_{\mathbb{R}^n}, \quad (\text{A.4})$$

and similarly for the side with $c_{\varphi,Q}$. Also, trivially, if $\|\varphi_*v\|_{\mathbb{R}^n} = 0$, then the same inequalities hold, proving (A.2).

Let us now turn to (A.3). Define

$$R_{\varphi,Q} = \{\xi \in \mathbb{R}^n \mid \exists (x, y, t) \in Q \times Q \times [0, 1], \xi = (1 - t)\varphi(x) + t\varphi(y)\},$$

which we note is compact. Let $x, y \in Q$ and define γ by

$$\gamma(t) = \varphi^{-1}((1 - t)\varphi(x) + t\varphi(y)), \quad \text{for all } t \in [0, 1],$$

so that in particular, the composition $\varphi \circ \gamma$ is the (Euclidean) geodesic from $\varphi(x)$ to $\varphi(y)$. Since $\varphi(V)$ is convex we deduce that $R_{\varphi,Q} \subset \varphi(V)$ and so $\gamma([0, 1]) \subset \varphi^{-1}(R_{\varphi,Q}) \subset U$. Since $\varphi^{-1}(R_{\varphi,Q})$ is furthermore compact, using (A.4) and setting $L_{\varphi,Q} = C_{\varphi, \varphi^{-1}(R_{\varphi,Q})}$ yields

$$\begin{aligned} d(x, y) &\leq \int_0^1 \|\gamma'(t)\|_{\gamma(t)} dt \leq L_{\varphi,Q} \int_0^1 \|\varphi_*\gamma'(t)\|_{\mathbb{R}^n} dt \\ &= L_{\varphi,Q} \int_0^1 \|(\varphi \circ \gamma)'(t)\|_{\mathbb{R}^n} dt = L_{\varphi,Q} \|\varphi(x) - \varphi(y)\|_{\mathbb{R}^n}. \end{aligned}$$

Appendix B. Proofs of preliminary results.

B.1. **Proof of Lemma 2.2.** Let us first show the following lemma, which is also going to be useful in Appendix B.2.

Lemma B.1. *Let $U \subset M$ be open and geodesically convex, and let further X be a differentiable vector field on U . Then,*

$$\|X(x) - \Pi_{yx}X(y)\|_x^2 \leq \int_0^1 \|\nabla_{\gamma'(t)}X(\gamma(t))\|_{\gamma(t)}^2 dt, \quad \text{for all } x, y \in U,$$

where $\gamma: [0, 1] \rightarrow M$ is the minimizing geodesic joining x to y . (Refer to (2.3) for the definition of the above integrand.)

Proof. Take two points $x, y \in U$ and the unique minimizing geodesic $\gamma: [0, 1] \rightarrow M$ joining the two points; in particular, $\gamma([0, 1]) \subset U$ for all $t \in [0, 1]$. Consider an orthonormal parallel frame $t \mapsto \{e_1(t), e_2(t), \dots, e_n(t)\}$ along γ . From

$$X(\gamma(t)) = \sum_{i=1}^n \langle X(\gamma(t)), e_i(t) \rangle_{\gamma(t)} e_i(t), \quad \text{for all } t \in [0, 1],$$

using that the parallel transport is a linear operator, we get:

$$\Pi_{yx} X(\gamma(1)) = \sum_{i=1}^n \langle X(\gamma(1)), e_i(1) \rangle_{\gamma(1)} \Pi_{yx} e_i(1) = \sum_{i=1}^n \langle X(\gamma(1)), e_i(1) \rangle_{\gamma(1)} e_i(0).$$

Consequently,

$$\|X(x) - \Pi_{yx} X(y)\|_x^2 = \sum_{i=1}^n (\langle X(\gamma(1)), e_i(1) \rangle_{\gamma(1)} - \langle X(\gamma(0)), e_i(0) \rangle_{\gamma(0)})^2. \quad (\text{B.1})$$

Using the compatibility of the connection with the metric [13, 14] we obtain

$$\begin{aligned} \frac{d}{dt} \langle X(\gamma(t)), e_i(t) \rangle_{\gamma(t)} &= \langle \nabla_{\gamma'(t)} X(\gamma(t)), e_i(t) \rangle_{\gamma(t)} + \langle X(\gamma(t)), \nabla_{\gamma'(t)} e_i(t) \rangle_{\gamma(t)} \\ &= \langle \nabla_{\gamma'(t)} X(\gamma(t)), e_i(t) \rangle_{\gamma(t)}, \end{aligned}$$

where for the second equality we used $\nabla_{\gamma'(t)} e_i(t) = 0$, for all $t \in [0, 1]$. By integrating the equation above from 0 to 1 and using it in (B.1) we find

$$\begin{aligned} \|X(x) - \Pi_{yx} X(y)\|_x^2 &= \sum_{i=1}^n \left(\int_0^1 \langle \nabla_{\gamma'(t)} X(\gamma(t)), e_i(t) \rangle_{\gamma(t)} dt \right)^2 \\ &\leq \int_0^1 \sum_{i=1}^n \langle \nabla_{\gamma'(t)} X(\gamma(t)), e_i(t) \rangle_{\gamma(t)}^2 dt \\ &= \int_0^1 \|\nabla_{\gamma'(t)} X(\gamma(t))\|_{\gamma(t)}^2 dt, \end{aligned}$$

where for the second line we used Jensen’s inequality, and for the third we used the fact that $\{e_i(t)\}_{i=1}^n$ is an orthonormal basis of $T_{\gamma(t)}M$. \square

Let us now turn to the proof of Lemma 2.2. Take $Q \subset B_\delta(x) \cap V$ compact and let $y, z \in Q$. We write $G_\varphi(y)$ for the metric tensor associated with $\langle \cdot, \cdot \rangle_y$ in the basis $\{\partial/\partial\varphi^1(y), \dots, \partial/\partial\varphi^n(y)\}$. For all $i, j \in \{1, \dots, n\}$ we denote by $p_i(y, z)$ and $f_{ij}(y, z)$ the i th coordinates in this basis of $X(y) - \Pi_{zy} X(z)$ and $\partial/\partial\varphi^j(y) - \Pi_{zy} \partial/\partial\varphi^j(z)$, respectively.

For all $i \in \{1, \dots, n\}$ we have

$$\begin{aligned} |p_i(y, z)| &= \left| \left\langle X(y) - \Pi_{zy} X(z), G_\varphi(y)^{-1} \frac{\partial}{\partial\varphi^i}(y) \right\rangle_y \right| \\ &\leq \|X(y) - \Pi_{zy} X(z)\|_y \left\| G_\varphi(y)^{-1} \frac{\partial}{\partial\varphi^i}(y) \right\|_y \leq P_\varphi \|X\|_{\text{Lip}(Q)} d(y, z), \end{aligned}$$

where $P_\varphi = \sup_{i \in \{1, \dots, n\}} \|G_\varphi^{-1}(\cdot) \partial/\partial\varphi^i(\cdot)\|_{L^\infty(Q)}$, which is finite since $\partial/\partial\varphi^i$ for all $i \in \{1, \dots, n\}$ and G_φ are continuous over the compact set Q ; note that P_φ is independent of y and z .

Similarly, for all $i, j \in \{1, \dots, n\}$, we have

$$|f_{ij}(y, z)| \leq P_\varphi F_\varphi d(y, z),$$

where $F_\varphi = \sup_{i \in \{1, \dots, n\}} \|\partial/\partial\varphi^i\|_{\text{Lip}(Q)}$, which is finite since $\partial/\partial\varphi^i$ as a vector field on $B_\delta(x)$ is Lipschitz continuous for all $i \in \{1, \dots, n\}$. Indeed, by Lemma B.1, for all $i \in \{1, \dots, n\}$ there holds

$$\left\| \frac{\partial}{\partial\varphi^i}(y) - \Pi_{zy} \frac{\partial}{\partial\varphi^i}(z) \right\|_y^2 \leq \int_0^1 \left\| \nabla_{\gamma'(t)} \frac{\partial}{\partial\varphi^i}(\gamma(t)) \right\|_{\gamma(t)}^2 dt, \quad \text{for all } y, z \in B_\delta(x),$$

where $\gamma: [0, 1] \rightarrow M$ is the unique minimizing geodesic linking y to z . Decomposing $\nabla_{\gamma'(t)} \frac{\partial}{\partial\varphi^i}(\gamma(t))$ using Christoffel symbols and using that the terms involved are continuous, hence locally bounded, one can bound the integrand in the right-hand side by $L \|\gamma'(t)\|_{\gamma(t)}$, with $L > 0$ constant. Then, by rescaling γ to be constant-speed (i.e., so that $\|\gamma'(t)\|_{\gamma(t)} = d(y, z)$ for all $t \in [0, 1]$), we get the desired Lipschitz continuity:

$$\left\| \frac{\partial}{\partial\varphi^i}(y) - \Pi_{zy} \frac{\partial}{\partial\varphi^i}(z) \right\|_y^2 \leq L \int_0^1 \|\gamma'(t)\|_{\gamma(t)}^2 dt = L d(y, z)^2, \quad \text{for all } y, z \in B_\delta(x).$$

Now, write $(x_1(y), \dots, x_n(y)) = \varphi_* X(y) \in \mathbb{R}^n$, so that

$$X(y) = x_1(y) \frac{\partial}{\partial\varphi^1}(y) + \dots + x_n(y) \frac{\partial}{\partial\varphi^n}(y),$$

and analogously for z (cf. Appendix A). By linearity of the parallel transport from z to y , we have

$$\begin{aligned} \sum_{i=1}^n p_i(y, z) \frac{\partial}{\partial\varphi^i}(y) &= X(y) - \Pi_{zy} X(z) = \sum_{i=1}^n x_i(y) \frac{\partial}{\partial\varphi^i}(y) - \sum_{i=1}^n x_i(z) \Pi_{zy} \frac{\partial}{\partial\varphi^i}(z) \\ &= \sum_{i=1}^n (x_i(y) - x_i(z)) \frac{\partial}{\partial\varphi^i}(y) + \sum_{i=1}^n x_i(z) \left(\frac{\partial}{\partial\varphi^i}(y) - \Pi_{zy} \frac{\partial}{\partial\varphi^i}(z) \right) \\ &= \sum_{i=1}^n (x_i(y) - x_i(z)) \frac{\partial}{\partial\varphi^i}(y) + \sum_{i=1}^n \sum_{j=1}^n x_j(z) f_{ij}(y, z) \frac{\partial}{\partial\varphi^i}(y), \end{aligned}$$

which yields, for all $i \in \{1, \dots, n\}$,

$$p_i(y, z) = x_i(y) - x_i(z) + \sum_{j=1}^n x_j(z) f_{ij}(y, z).$$

Thus, for all $i \in \{1, \dots, n\}$,

$$\begin{aligned} |x_i(y) - x_i(z)| &\leq |p_i(y, z)| + \sum_{j=1}^n |x_j(z)| |f_{ij}(y, z)| \\ &\leq P_\varphi (\|X\|_{\text{Lip}(Q)} + n P_\varphi F_\varphi \|X\|_{L^\infty(Q)}) d(y, z). \end{aligned}$$

To conclude the proof we use (A.3). Indeed, by definition of normal chart, we have $\varphi(B_\delta(x)) = \varphi(\exp_x(B_\delta(0))) = E(\log_x(\exp_x(B_\delta(0)))) = E(B_\delta(0)) \cong B_\delta(0)$, which is a convex subset of \mathbb{R}^n , where E is given in (A.1).

B.2. Proof of Lemma 2.3. Assume first that ∇f is locally Lipschitz continuous. Let $Q \subset U$ be compact, take any $x \in Q$ and $v \in T_x M$ and write γ for the constant-speed geodesic with $\gamma(0) = x$ and $\gamma'(0) = v$. In particular, $d(\gamma(t), x) = t\|v\|_x$ for all $t \geq 0$ such that $\gamma(t)$ is defined, and for $\tau > 0$ small enough we have $\gamma(t) \in \tilde{Q}$ for all $t < \tau$ for some \tilde{Q} compact satisfying $Q \subset \tilde{Q} \subset U$. By local Lipschitz continuity of ∇f we thus have, for all $t < \tau$,

$$\|\Pi_{x\gamma(t)}^{-1} \nabla f(\gamma(t)) - \nabla f(x)\|_x \leq L_{\tilde{Q}} d(\gamma(t), x) = L_{\tilde{Q}} t\|v\|_x. \tag{B.2}$$

By (2.3) and the definition of the Hessian, we have

$$\text{Hess}_v f(x) = \nabla_v(\nabla f)(x) = \lim_{t \rightarrow 0} \frac{\Pi_{x\gamma(t)}^{-1} \nabla f(\gamma(t)) - \nabla f(x)}{t},$$

and using (B.2) we find

$$\|\text{Hess}_v f(x)\|_x \leq L_{\tilde{Q}} \|v\|_x.$$

The Hessian operator $\text{Hess} f(x)$ is therefore bounded. By the arbitrariness of Q and $x \in Q$, we conclude that $\text{Hess} f$ is locally bounded.

To prove the converse, suppose that U is convex and that the Hessian of f is locally bounded. Let $Q \subset U$ be compact; because U is open and convex, there exists a compact, geodesically convex set \tilde{Q} with $Q \subset \tilde{Q} \subset U$. Furthermore, take two points $x, y \in Q$ and the constant-speed geodesic $\gamma: [0, 1] \rightarrow M$ joining the two points; in particular, $\|\gamma'(t)\|_{\gamma(t)} = d(x, y)$ and $\gamma([0, 1]) \subset \tilde{Q}$ for all $t \in [0, 1]$. By Lemma B.1 applied to the vector field ∇f we get

$$\|\nabla f(x) - \Pi_{yx} \nabla f(y)\|_x^2 \leq \int_0^1 \|\text{Hess}_{\gamma'(t)} f(\gamma(t))\|_{\gamma(t)}^2 dt.$$

By the local boundedness of $\text{Hess} f$ and the fact that γ is constant-speed we further get

$$\|\nabla f(x) - \Pi_{yx} \nabla f(y)\|_x^2 \leq C_Q^2 d(x, y)^2,$$

for some $C_Q > 0$, which, by arbitrariness of Q and $x, y \in Q$, ends the proof.

B.3. Proof of Theorem 2.5. Take $x \in U$ and let $\delta \leq r_{\text{conv}}(x)$ be such that $B_\delta(x) \subset U$. Let $(B_\delta(x), \varphi)$ be a normal chart generated by x , and consider the initial-value problem

$$\begin{cases} \alpha'(t) = \Xi(\alpha(t), t), \\ \alpha(0) = \varphi(x), \end{cases} \tag{B.3}$$

where we define $\Xi: \varphi(B_\delta(x)) \times [0, T] \rightarrow \mathbb{R}^n$ by

$$\Xi(\xi, t) = \varphi_* X_t(\varphi^{-1}(\xi)), \quad \text{for all } (\xi, t) \in \varphi(B_\delta(x)) \times [0, T].$$

Take $R \subset \varphi(B_\delta(x)) = B_\delta(0)$ and $S \subset [0, T]$ compact, so that in particular $Q := \varphi^{-1}(R) \subset B_\delta(x)$ is compact. For all $\xi, \eta \in R$ and $t \in S$, our Lipschitz-continuity assumption on X_t yields

$$\begin{aligned} \|\Xi(\xi, t) - \Xi(\eta, t)\|_{\mathbb{R}^n} &= \|\varphi_* X_t(\varphi^{-1}(\xi)) - \varphi_* X_t(\varphi^{-1}(\eta))\|_{\mathbb{R}^n} \\ &\leq (A \|X_t\|_{L^\infty(Q)} + B \|X_t\|_{\text{Lip}(Q)}) \|\xi - \eta\|_{\mathbb{R}^n}, \end{aligned}$$

for some constants $A, B > 0$ coming from Lemma 2.2. Also, for all $\xi \in R$ it holds that

$$\|\Xi(\xi, t)\|_{\mathbb{R}^n} = \|\varphi_* X_t(\varphi^{-1}(\xi))\|_{\mathbb{R}^n} \leq C \|X_t\|_{L^\infty(Q)},$$

where $C > 0$ comes from (A.2). Therefore, by (2.5) we get that Ξ satisfies

$$\int_S \left(\|\Xi(\cdot, t)\|_{L^\infty(R)} + \|\Xi(\cdot, t)\|_{\text{Lip}(R)} \right) dt \leq \int_S \left((A + C) \|X_t\|_{L^\infty(Q)} + B \|X_t\|_{\text{Lip}(Q)} \right) dt < \infty.$$

By arbitrariness of the compact sets $R \subset \varphi(B_\delta(x))$ and $S \subset [0, T)$ and by the classical Cauchy–Lipschitz theorem on \mathbb{R}^n (see for instance [3, Lemma 8.1.4]), this yields the existence of a unique maximal solution α_x to (B.3) defined on some time interval $[0, \tau_x)$, with $\tau_x \leq T$, and with values in $\varphi(B_\delta(x))$. By defining $\varphi_x = \varphi^{-1} \circ \alpha_x$, we see that α_x satisfies (B.3) if and only if

$$\begin{cases} \varphi_* \varphi'_x(t) = (\varphi \circ \varphi_x)'(t) = \varphi_* X_t(\varphi_x(t)) & \text{for all } t \in [0, \tau_x), \\ \varphi(\varphi_x(0)) = \varphi(x). \end{cases}$$

By the bijectivity of φ , we get that φ_x is thus the unique maximal solution to the characteristic equation starting at x .

Let now Σ be a compact subset of U . We are left with showing that $\tau := \inf_{x \in \Sigma}(\tau_x) > 0$. By classical Euclidean Lipschitz theory, we deduce that for all $x \in U$ there exists $\delta_x > 0$ such that $\bar{\tau}_x := \inf_{y \in B_{\delta_x}(x)} \tau_y > 0$. Since Σ is compact we know it can be covered by a finite subfamily of $\{B_{\delta_x}(x)\}_{x \in \Sigma}$, which we index by $\{x_1, \dots, x_n\}$ for some $n \in \mathbb{N}$. We thus get $\tau = \min_{i \in \{1, \dots, n\}} \bar{\tau}_{x_i} > 0$. The map Ψ defined by $\Psi(x, t) = \varphi_x(t)$ for all $x \in \Sigma$ and $t \in [0, \tau)$ is then the unique maximal flow map generated by (X, Σ) .

B.4. Proof of Theorem 2.6. We distinguish two cases: M is compact, and M is unbounded.

The first case is direct by the Escape Lemma (cf. [31, Exercise 4.10] for instance), which states that local solutions cannot be contained in any compact subset of M . The local Cauchy–Lipschitz theorem (cf. Theorem 2.5) together with a contradiction argument suffices then to end this case.

Suppose now that M is unbounded. First note that the bound assumption on the Lipschitz constants of our vector fields implies that there exists $C > 0$ such that for any $U \subset M$ we have

$$\sup_{t \in [0, \infty)} \|X_t\|_{L^\infty(U)} \leq C(1 + \text{diam}(U)). \tag{B.4}$$

Let $\Sigma \subset M$ be compact and U be bounded and open and such that $\Sigma \subset U$. Since M is convex we know U is a totally normal neighborhood. Denote by \tilde{X} the restriction of X to $U \times [0, \infty)$. By the local Cauchy–Lipschitz theorem (cf. Theorem 2.5), there exists a unique maximal flow map generated by (\tilde{X}, Σ) associated with some maximal time of existence $\tau > 0$. Recalling the proof of Theorem 2.5 and the classical Cauchy–Lipschitz theorem in Euclidean setting, we know there exists a constant $\alpha > 0$ such that

$$\tau \geq \alpha \min \left(\frac{\text{dist}(\Sigma, \partial U)}{\sup_{t \in [0, \infty)} \|X_t\|_{L^\infty(U)}}, \frac{1}{\sup_{t \in [0, \infty)} \|X_t\|_{\text{Lip}(M)}} \right),$$

where ∂U is the boundary of U . By (B.4) we then have

$$\tau \geq \alpha \min \left(\frac{\text{dist}(\Sigma, \partial U)}{C(1 + \text{diam}(U))}, \frac{1}{\sup_{t \in [0, \infty)} \|X_t\|_{\text{Lip}(M)}} \right).$$

Because M is unbounded, we see that we can always choose U a posteriori so that $\frac{\text{dist}(\Sigma, \partial U)}{1 + \text{diam}(U)} \geq 1$. Overall, this makes our lower bound on the maximal time of existence independent of Σ and U . Applying a simple iterative argument extends our restricted local flow map to a unique global one, thus showing the desired result.

B.5. Proof of Lemma 4.2. The proof follows from (4.7). Fix $z \in U$, and note first that

$$U \subset B_r(z) \subset B_{2r}(p).$$

Both inclusions follow from the triangle inequality, using that $U \subset B_{r/2}(p)$. Indeed, for $x \in U \subset B_{r/2}(p)$, one has $d(x, p) < r/2$, and hence, $d(x, z) \leq d(x, p) + d(p, z) < r/2 + r/2 = r$, showing the first inclusion. For the second inclusion, take $y \in B_r(z)$ and estimate $d(y, p) \leq d(y, z) + d(z, p) < r + r/2 < 2r$.

In the set $B_{2r}(p)$, the sectional curvature is bounded below by λ and above by μ , so that (4.6) holds on $B_r(z)$. Consequently, (4.7) holds for all $x, z \in U \subset B_r(z)$. By the assumption (4.8) on the diameter of U , the left-hand side in (4.7) is nonnegative (when $\mu > 0$, use $\sqrt{\mu} d(x, z) \leq \sqrt{\mu} \Delta < \frac{\pi}{2}$). On the other hand, for the right-hand side of (4.7), use that the function $s \mapsto s \coth(s)$ is nondecreasing, and so

$$\sqrt{-\lambda} d(x, z) \coth(\sqrt{-\lambda} d(x, z)) \leq \sqrt{-\lambda} \Delta \coth(\sqrt{-\lambda} \Delta) = \frac{L}{2}, \quad \text{for all } x, z \in U.$$

From these observations, (4.7) yields

$$0 \leq \langle \text{Hess}_v d_z^2(x), v \rangle_x \leq L \|v\|_x^2, \quad \text{for all } x, z \in U \text{ and } v \in T_x M. \tag{B.5}$$

Since the Hessian operator is symmetric, for any $x \in U$ one can choose an orthonormal basis of $T_x M$ consisting of eigenvectors of the Hessian of d_z^2 at x . By decomposing into this basis and using (B.5) for the elements of the basis, one can get, via the Cauchy–Schwartz inequality,

$$|\langle \text{Hess}_v d_z^2(x), w \rangle_x| \leq L \|v\|_x \|w\|_x, \quad \text{for all } x, z \in U \text{ and } v, w \in T_x M.$$

Consequently, we can bound the operator (and hence, the Euclidean) norm of $\text{Hess}_v d_z^2(x)$ to get

$$\| \text{Hess}_v d_z^2(x) \|_x \leq L \|v\|_x, \quad \text{for all } x, z \in U \text{ and } v, w \in T_x M.$$

Finally, if U is convex, then (4.9) follows directly from Lemma 2.4 and (4.2).

B.6. Proof of Lemma 4.3. For completeness, let us give the statement of [29, Corollary 6.6.1], which lies at the core of the proof of Lemma 4.3. Given a differentiable function $f: M_1 \rightarrow M_2$ between two smooth Riemannian manifolds M_1 and M_2 , we write $D_x f: T_x M_1 \rightarrow T_{f(x)} M_2$ for the differential map of f at x .

Lemma B.2. *Let $p \in M$ and $r < \text{inj}(M)$. Let there be a minimizing geodesic connecting two distinct points x and y in $B_r(p)$, with velocity u at the point x . Write $\mu = \mu_r(p)$, and suppose that $\|u\|_x \leq \frac{\pi}{\sqrt{\mu}}$ in case $\mu > 0$, and let $w \in T_x M$ be so that $\langle w, u \rangle_x = 0$. Then,*

$$\|w\|_x \frac{s_\mu(\|u\|_x)}{\|u\|_x} \leq \|D_u \exp_x(w)\|_y, \tag{B.6}$$

where s_μ is defined, for all $a \geq 0$, by

$$s_\mu(a) = \begin{cases} \frac{\sin(a\sqrt{\mu})}{\sqrt{\mu}} & \text{if } \mu > 0, \\ a & \text{if } \mu = 0. \end{cases}$$

Let us now prove Lemma 4.3. The case $\mu = 0$ is stated in [29, Corollary 6.9.1]. We give its proof here for completeness. Fix $x, y \in U$ with $x \neq y$. Because U is a totally normal neighborhood there exists only one minimizing geodesic linking x to y with velocity $\log_x y$ at the origin, and the exponential map \exp_x is a diffeomorphism from $T_x M$ to U . Then, we can let $v \in T_y M$ and apply Lemma B.2 with $w = D_y \log_x(v)$ and $u = \log_x y$. Note that since $\log_x : M \rightarrow T_x M$, we have $D_y \log_x(v) \in T_{\log_x y} T_x M \cong T_x M$. Also, up to a rotation, we can assume $\langle \log_x y, D_y \log_x(v) \rangle_x = 0$. By (B.6), we then get

$$\|w\|_x \leq \|D_{\log_x y} \exp_x(w)\|_y = \|D_{\log_x y} \exp_x(D_y \log_x(v))\|_y = \|v\|_y,$$

where for the last equality we used $\exp_x \circ \log_x$ is the identity map on U . We give the obtained result as:

$$\|D_y \log_x(v)\|_x \leq \|v\|_y, \quad \text{for all } x, y \in M, v \in T_y M. \tag{B.7}$$

Now fix $x, y, z \in U$, and denote by γ the constant-speed geodesic between y and z . Then,

$$\log_x z - \log_x y = \int_0^1 \frac{d}{dt} \log_x \gamma(t) dt = \int_0^1 D_{\gamma(t)} \log_x(\gamma'(t)) dt. \tag{B.8}$$

By applying (B.7) in the integral above for $y = \gamma(t)$ and $v = \gamma'(t)$, we get from (B.8):

$$\|\log_x z - \log_x y\|_x \leq \int_0^1 \|D_{\gamma(t)} \log_x(\gamma'(t))\|_x dt \leq \int_0^1 \|\gamma'(t)\|_{\gamma(t)} dt, \tag{B.9}$$

which leads to (4.10) given that $\|\gamma'(t)\| = d(y, z)$ for all $t \in [0, 1]$.

Let us turn to the proof for the case $\mu > 0$, which is again based on Lemma B.2. Fix $x, y \in U$ with $x \neq y$. Similar to the $\mu = 0$ case, there exists only one minimizing geodesic joining x and y with velocity $\log_x y$ at x , and the exponential map \exp_x is a diffeomorphism from $T_x M$ to U . By the assumption on the diameter of U we have $d(x, y) \leq \frac{\pi - \varepsilon}{\sqrt{\mu}} < \frac{\pi}{\sqrt{\mu}}$, and hence we can then let $v \in T_y M$ and apply (B.6) for $w = D_y \log_x(v)$ and $u = \log_x y$, to get

$$\|w\|_x \leq \frac{\sqrt{\mu} \|\log_x y\|_x}{\sin(\sqrt{\mu} \|\log_x y\|_x)} \|D_{\log_x y} \exp_x(w)\|_y = \frac{\sqrt{\mu} \|\log_x y\|_x}{\sin(\sqrt{\mu} \|\log_x y\|_x)} \|v\|_y,$$

and hence,

$$\|D_y \log_x(v)\|_x \leq \frac{\sqrt{\mu} \|\log_x y\|}{\sin(\sqrt{\mu} \|\log_x y\|)} \|v\|_y.$$

Since the function $\tau \mapsto \frac{\tau}{\sin \tau}$ is nondecreasing on $[0, \pi - \varepsilon]$, we get from above that

$$\|D_y \log_x(v)\|_x \leq \frac{\pi - \varepsilon}{\sin(\pi - \varepsilon)} \|v\|_y, \quad \text{for all } x, y \in U, \text{ and } v \in T_y M. \tag{B.10}$$

Now fix $x, y, z \in U$ and take γ the constant-speed geodesic between y and z . Restarting from (B.8), equation (B.10) yields (see also (B.9))

$$\|\log_x z - \log_x y\|_x \leq \frac{\pi - \varepsilon}{\sin(\pi - \varepsilon)} \int_0^1 \|\gamma'(t)\|_{\gamma(t)} dt.$$

The Lipschitz estimate (4.11) now follows from above, using that $\|\gamma'(t)\| = d(y, z)$ for all $t \in [0, 1]$. Note that the Lipschitz constant approaches ∞ as $\varepsilon \rightarrow 0$.

Acknowledgements. The authors thank Xavier Pennec for his time and insightful discussions on the topic of bounded curvature.

REFERENCES

- [1] F. Alimisis, A. Orvieto, G. Bécigneul and A. Lucchi, A continuous-time perspective for modeling acceleration in riemannian optimization, *Int. Conf. Art. Intel. Statist.*, (2020), 1297–1307.
- [2] Foivos Alimisis, Antonio Orvieto, Gary Bécigneul and Aurelien Lucchi, Momentum improves optimization on Riemannian manifolds, [arXiv:2002.04144](https://arxiv.org/abs/2002.04144), 2021.
- [3] L. Ambrosio, N. Gigli and G. Savaré, *Gradient Flows in Metric Spaces and in the Space of Probability Measures*, Lectures in Mathematics ETH Zürich. Birkhäuser Verlag, Basel, 2005.
- [4] Andrea L. Bertozzi, José A. Carrillo and Thomas Laurent, Blow-up in multidimensional aggregation equations with mildly singular interaction kernels, *Nonlinearity*, **22** (2009), 683–710.
- [5] Andrea L. Bertozzi and Thomas Laurent, Finite-time blow-up of solutions of an aggregation equation in \mathbb{R}^n , *Comm. Math. Phys.*, **274** (2007), 717–735.
- [6] Andrea L. Bertozzi, Thomas Laurent and Jesús Rosado, L^p theory for the multidimensional aggregation equation, *Commun. Pure Appl. Math.*, **64** (2011), 45–83.
- [7] M. Bodnar and J. J. L. Velazquez, An integro-differential equation arising as a limit of individual cell-based models, *J. Differ. Equ.*, **222** (2006), 341–380.
- [8] J. A. Carrillo, M. Di Francesco, A. Figalli, T. Laurent and D. Slepčev, Global-in-time weak measure solutions and finite-time aggregation for nonlocal interaction equations, *Duke Math. J.*, **156** (2011), 229–271.
- [9] J. A. Carrillo, D. Slepčev and L. Wu, Nonlocal-interaction equations on uniformly prox-regular sets, *Discrete Contin. Dyn. Syst. Ser. A*, **36** (2016), 1209–1247.
- [10] José A. Carrillo, Robert J. McCann and Cédric Villani, Contractions in the 2-Wasserstein length space and thermalization of granular media, *Arch. Ration. Mech. Anal.*, **179** (2006), 217–263.
- [11] I. Chavel, *Riemannian Geometry : A Modern Introduction*, Cambridge Studies in Advanced Mathematics. Cambridge University Press, second edition, 2006.
- [12] Fabiana R. de Oliveira and Orizon Pereira Ferreira, Newton method for finding a singularity of a special class of locally Lipschitz continuous vector fields on Riemannian manifolds, *J. Optim. Theory Appl.*, 185 (2020), 522–539.
- [13] M. P. do Carmo, *Riemannian Geometry*, Mathematics: Theory and Applications. Birkhäuser, Boston, second edition, 1992.
- [14] A. Durmus, P. Jiménez, E. Moulines, S. Said and H.-T. Wai, Convergence analysis of Riemannian stochastic approximation schemes, [arXiv:2005.13284](https://arxiv.org/abs/2005.13284), 2020.
- [15] K. Fellner and G. Raoul, Stable stationary states of non-local interaction equations, *Math. Models Methods Appl. Sci.*, 20 (2010), 2267–2291.
- [16] O. P. Ferreira, M. S. Louzeiro and L. F. Prudente, Gradient method for optimization on Riemannian manifolds with lower bounded curvature, *SIAM J. Optim.*, 29 (2019), 2517–2541.
- [17] R. C. Fetecau, S-Y. Ha and H. Park, An intrinsic aggregation model on the special orthogonal group $SO(3)$: well-posedness and collective behaviours, *J. Nonlinear Sci.*, 31 (2021), 74.
- [18] R. C. Fetecau and Y. Huang, Equilibria of biological aggregations with nonlocal repulsive-attractive interactions, *Phys. D*, **260** (2013), 9–64.
- [19] R. C. Fetecau, Y. Huang and T. Kolokolnikov, Swarm dynamics and equilibria for a nonlocal aggregation model, *Nonlinearity*, **24** (2011), 2681–2716.
- [20] R. C. Fetecau, H. Park and F. S. Pataccini, Well-posedness and asymptotic behaviour of an aggregation model with intrinsic interactions on sphere and other manifolds, *Anal. Appl.*, **19** (2021), 965–1017.
- [21] R. C. Fetecau and B. Zhang, Self-organization on Riemannian manifolds, *J. Geom. Mech.*, **11** (2019), 397–426.
- [22] V. Gazi and K. M. Passino, Stability analysis of swarms, In *Proc. American Control Conf.*, pages 8–10, Anchorage, AK, 2002.
- [23] S-Y. Ha and H. Park, Emergent behaviors of Lohe tensor flock, *J. Stat. Phys.*, **178** (2020), 1268–1292.
- [24] Seung-Yeal Ha, Dongnam Ko and Sang Woo Ryoo, Emergent dynamics of a generalized Lohe model on some class of Lie groups, *J. Stat. Phys.*, **168** (2017), 171–207.

- [25] Seung-Yeal Ha, Dongnam Ko and Sang Woo Ryoo, [On the relaxation dynamics of Lohe oscillators on some Riemannian manifolds](#), *J. Stat. Phys.*, **172** (2018), 1427–1478.
- [26] Zhang Hongyi and Suvrit Sra, [Towards Riemannian accelerated gradient methods](#), [arXiv:1806.02812](#), 2018.
- [27] M. C. Irwin, *Smooth Dynamical Systems*, volume 17 of *Advanced Series in Nonlinear Dynamics*, World Scientific, 2001.
- [28] M. Ji and M. Egerstedt, Distributed coordination control of multi-agent systems while preserving connectedness, *IEEE Trans. Robot.*, **23** (2007), 693–703.
- [29] J. Jost, *Riemannian Geometry and Geometric Analysis*, Universitext. Springer-Verlag, seventh edition, 2017.
- [30] Theodore Kolokolnikov, Hui Sun, David Uminsky and Andrea L. Bertozzi, A theory of complex patterns arising from 2D particle interactions, *Phys. Rev. E, Rapid Communications*, **84** (2011), 015203(R), 2011.
- [31] J. M. Lee, *Introduction to Riemannian manifolds*, volume 176 of *Graduate Texts in Mathematics*, Springer International Publishing, second edition, 2018.
- [32] Andrew J. Leverentz, Chad M. Topaz and Andrew J. Bernoff, [Asymptotic dynamics of attractive-repulsive swarms](#), *SIAM J. Appl. Dyn. Syst.*, **8** (2009), 880–908.
- [33] A. Mihai, Affine transformations on Banach manifolds, *Atti della Accademia Nazionale dei Lincei. Classe di Scienze Fisiche, Matematiche e Naturali. Rendiconti*, **8** (1976), 224–227.
- [34] A. Mogilner and L. Edelstein-Keshet, [A non-local model for a swarm](#), *J. Math. Biol.*, **38** (1999), 534–570.
- [35] Sebastien Motsch and Eitan Tadmor, [Heterophilious dynamics enhances consensus](#), *SIAM Review*, **56** (2014), 577–621.
- [36] F. S. Patacchini and D Slepčev, [The nonlocal-interaction equation near attracting manifolds](#), [arXiv:2106.01823](#), 2021.
- [37] James von Brecht, David Uminsky, Theodore Kolokolnikov and Andrea Bertozzi, [Predicting pattern formation in particle interactions](#), *Math. Models Methods Appl. Sci.*, **22** (2012), (Supp. 1):1140002.
- [38] L. Wu and D. Slepčev, [Nonlocal interaction equations in environments with heterogeneities and boundaries](#), *Commun. Partial Differ. Equ.*, **40** (2015), 1241–1281.

Received September 2021; revised June 2022; early access July 2022.

E-mail address: van@math.sfu.ca

E-mail address: francesco.patacchini@ifpen.fr